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Spatial scaling metrics of mask-induced induced line-edge roughness

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Abstract

Mask contributors to line-edge roughness (LER) have recently been shown to be an issue of concern for both the accuracy of current resist evaluation tests as well the ultimate LER requirements for the 22-nm production node. Problems arise from mask absorber LER as well as mask reflector or surface roughness leading to random phase variations in the reflected beam. Not only do these mask contributors effect to total measured LER in resist, but they also effect LER spatial characteristic such as the LER power spectral density and related descriptors of correlation length and roughness exponent. Noting that these characteristics are important in the understanding of LER, it is crucial to understand how mask effects impact these parameters. Moreover, understanding how these metrics respond to mask effects may lead to an experimental mechanism for experimentally evaluating the importance of mask contributors to LER. Here we use computer modeling to study the LER spatial metrics arising from mask effects. We further describe the effects of illumination conditions and defocus on the metrics and compare the results to those expected from intrinsic resist LER.

Keyword: extreme ultraviolet, lithography, multilayer, mask, line-edge roughess

Introduction

Mask contributors to line-edge roughness (LER) have recently been shown [1] to be an issue of concern for both the accuracy of current resist evaluation tests as well the ultimate LER requirements for the 22-nm production node [2]. Problems arise from mask absorber LER [3] as well as mask reflector or surface roughness leading to random phase variations in the reflected beam [4, 5]. Not only do these mask contributors effect to total measured LER in resist, but they may also effect LER spatial characteristic such as correlation length and roughness exponent which have been proposed as important metrics for understanding LER [6-8]. Moreover, understanding these additional metrics and how they might vary compared to values expected from resist may provide an experimental mechanism for evaluating the importance of mask contributors to LER.

Here we explicitly consider the spatial metrics of mask-induced LER through computer modeling and study their potential impact on experimentally measured resist LER metrics.

Mask-induced LER spatial metrics

The modeling method used in this study is identical to that described in the literature [1] and is only briefly presented here. Thin mask modeling is used and the multilayer surface roughness is represented as a random phase error on the mask. The morphology of the absorber LER is set match absorber LER as measured using scanning electron micrographs (SEMs) and the phase roughness morphology is determined from atomic force microscopy [1]. For the optic characteristics in the modeling, we use published wavefront aberration and flare measurements from the SEMATECH Berkeley microfield exposure tool (MET) [9-12].

For reference, we begin by considering the expected spatial LER metrics for photoresist-induced LER. A stochastic resist model has recently been developed allowing the spatial characteristics of resist LER to be accurately predicted [13]. This model has been verified against

a variety of resist and demonstrated remarkable accuracy in the prediction of the shape of the LER power spectral density (PSD) [14] for PSD magnitudes above the scanning electron microscope noise floor. Figure 1 shows a summary of the comparison results for four different resist platforms. The reported R values correspond to the width of the deprotection blur function at the heart of the resist model [13] and is closely related to the correlation length of the LER. To the extent that the model captures the dominant mechanisms leading to LER, the R term extracted by the model can be seen to represent the intrinsic resolution limit of the resist.

Next we consider the LER spatial metrics on the mask itself. Figure 2 shows a SEM of 250-nm line-space patterns on the 5× mask corresponding to 50-nm patterns at the wafer. The mask LER PSD (scaled to wafer coordinates) is shown in Fig. 3. The mask LER is well fit by the stochastic resist LER model used above. In wafer coordinates, the roughness exponent of the LER is 0.83 and the correlation length is 11.4 nm with both values being determined from the height-height correlation function (HHCF) [7, 15]. In mask coordinates, the correlation length is simply 5 times larger while the roughness exponent remains the same.

Although the mask LER itself is very similar in morphology to that seen in imaged resist, it is important to recall that the morphology in the projected image LER will not be the same as the mask LER. LER transfer from the mask to the wafer is dominated by the LER Transfer Function (LTF) [3] which is distinct from the optic contrast transfer function. In practice the LTF is best determined through computer simulation readily allowing system parameters such as pupil aberrations and illumination pupil fill to be incorporated in the prediction [3]. In Fig. 4 we show the LER PSD of the projected mask image directly comparing it to the mask LER PSD from Fig. 3. The optical parameters are set to match those of the MET with annular 0.35-0.55 illumination. The effect of the LTF on the LER PSD is evident and clearly changes the

morphology of the PSD such that it is no longer well represented by a self-affine fractal roughness as predicted by the resist LER model and as observed in resists. This striking difference in the observed PSD suggests that the PSD is a useful method for ascertaining the importance of mask LER in printed LER. Although the effect of the LTF is clearly visible in the PSD, the height-height correlation and sigma versus length functions do not exhibit a significant difference in morphological signature as compared to self-affine roughness as observed from resists (Fig. 5). These functions, thus, do not serve as good root cause indicators for printed LER. After projection by the optic, the HHCF measured correlation length and roughness exponent are 32.6 nm and 0.90, respectively. Not much change in the roughness exponent but a significant increase in the correlation length owing the band-limiting characteristics of the LTF. The PSD-measured correlation length is 29.6 nm. The PSD-measured roughness exponent as determined from the slope [roughness exponent is $(Slope_{PSD}-1)/2$] is greater than 5 which is clearly not representative of a self-affine process.

Next we consider the LER spatial metrics for mask multilayer roughness induced LER. The same modeling conditions as described for the mask LER case are used except that the mask absorber pattern is now assumed to be perfect. The LER spatial metrics shown in Fig. 6 again show the mask roughness induced LER PSD to deviate strongly in morphology compared to self-affine roughness, whereas the HHCF and sigma versus length plots due not capture this change. The HHCF-measured correlation length and roughness exponent are 16.2 nm and 0.97, whereas the PSD-measured correlation length is 29.2 nm. Again, the PSD-measured roughness exponent as being greater than 6 is not representative of a self-affine process. The striking similarity between the PSD results for the two different types of mask roughness can be understood by realizing that the spatial characteristics of the projected multilayer roughness are dominated by

the transfer function of the projection optics. Moreover considering the transfer all the way to the LER instead of simply intensity variations at the wafer, it is evident that the characteristics of the closely related LTF will dominate as is the case for the more evident condition of mask pattern LER considered above.

These results indicate a potential method for ascertaining the importance of mask effects in measured LER: namely looking for large discrepancies between the roughness exponent measured using the HHCF and PSD methods, respectively. Also, a PSD roughness exponent measurement of greater than one may indicate some level of mask contribution to the measured LER.

The results above were all obtained assuming conventional annular illumination on the centrally obscured MET. It was previously shown, however, the illumination coherence settings play an important role in the magnitude of the LER, especially for the mask multilayer roughness case. Thus we next consider the impact of illumination settings on the spatial metrics. Table 1 shows the extracted correlation length and roughness exponent for annular and dipole illumination settings on the two mask types. The annular settings are as described above and the dipole settings are pole σ of 0.1, pole offset of 0.57, and an orientation of 45°. The results in Table 1 along with Figs. 7 and 8 show the illumination settings to have minimal impact on the extracted LER spatial metrics. The expected LER magnitude changes as reported in Ref. [1], however, are evident from the plots in Figs. 7 and 8. Dipole illumination improves the printed LER for the case of LER on the mask and degrades the printed LER for the case of mask multilayer roughness.

Focus settings were also found to have significant impact on the mask-induced LER magnitude, especially for the case of the mask multilayer roughness [1]. Finally we consider

these same effects but from the perspective of LER spatial metrics. Table 2 shows the LER metric results for the aerial image at best focus and 110-nm defocus with annular illumination. 110-nm was chosen because this corresponds to the edge of the focus latitude based on an aerial-image contrast limit of 50%. The results are inconclusive for the HHCF method which shows a possible increase in measured correlation length for the mask LER case and a possible decrease in roughness exponent for the mask multilayer roughness case. The PSD method, on the other hand, shows a significant and consistent reduction in measured slope as seen from the roughness exponent. Figures 9 and 10 shows the corresponding LER characterization plots comparing the two different focus settings. These results again point to the PSD metrics as a potential mechanism for ascertaining the importance of mask effects in printed LER, this time by observing changes through focus. The focus results in Figs. 9 and 10 also indicate that the PSD could be used to distinguish between the two different forms of mask-induced LER.

Discussion

The results presented here indicate that the spatial LER metrics as observed from the LER PSD differ significantly for intrinsic resist LER compared to mask-induced LER. The difference, however, is not evident in the height-height correlation and sigma versus length functions. These results show that the LER PSD could be useful in ascertaining the importance of mask effects in the observed LER. Moreover, through-focus observation of the LER PSD could even allow the effects of mask absorber LER to be distinguished from mask multilayer-roughness induced LER.

The fact that current resist LER typically exhibits the expected self-affine behavior in the PSD as observed in Fig. 1 suggests that mask roughness is presently not a dominant effect in observed LER limits which is consistent with previous conclusions [1].

Although the potential utility of using the LER PSD to ascertain the importance of mask-induced LER has been demonstrated, exactly how sensitive this approach may be in practice has not yet been determined. This question will be addressed in future studies where the stochastic resist model effects [13] will be explicitly added to the simulations.

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List of Figures

Fig. 1. A summary of comparison results between experimental resist LER and the stochastic resist model for four different resist platforms. The reported R values correspond to the width of the deprotection blur function at the heart of the resist model [13] and is closely related to the correlation length of the LER.

Fig. 2. SEM of 250-nm line-space patterns on the 5 mask corresponding to 50-nm patterns at the wafer.

Fig. 3. The mask LER PSD (scaled to wafer coordinates) compared to the stochastic resist LER model (dashed line).

Fig. 4. LER PSD of the projected mask image compared to the mask LER PSD from Fig. 3. A significant difference in morphology compared to the expected resist LER shape is observed.

Fig. 5. LER height-height correlation and sigma versus length functions for the projected mask image. No significant morphological differences are evident between the projected mask LER functions and those expected from the stochastic resist model.

Fig. 6. LER PSD, height-height correlation and sigma versus length functions for the projected mask image for the mask multilayer roughness case. Again, morphological differences between the mask-induced LER and expected resist LER are only observed in the PSD.

Fig. 7. Projected LER spatial metrics for annular and dipole illumination considering mask absorber LER alone. (a) PSD, (b) HHCF, (c) sigma versus length.

Fig. 8. Projected LER spatial metrics for annular and dipole illumination considering mask multilayer roughness alone. (a) PSD, (b) HHCF, (c) sigma versus length.

Fig. 9. Projected LER spatial metrics for annular illumination and two different focus settings (best focus and 110 nm defocus) considering mask absorber LER alone. (a) PSD, (b) HHCF, (c) sigma versus length.

Fig. 10. Projected LER spatial metrics for annular illumination and two different focus settings (best focus and 110 nm defocus) considering mask multilayer roughness alone. (a) PSD, (b) HHCF, (c) sigma versus length.
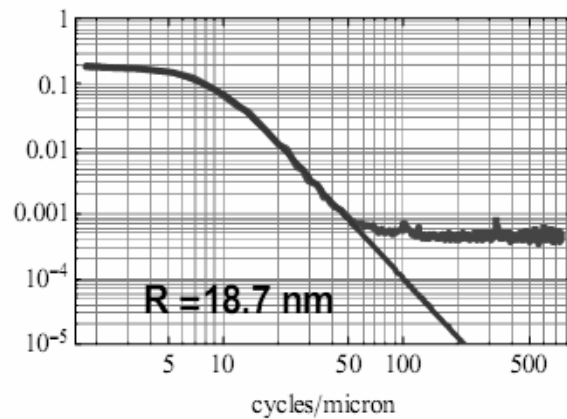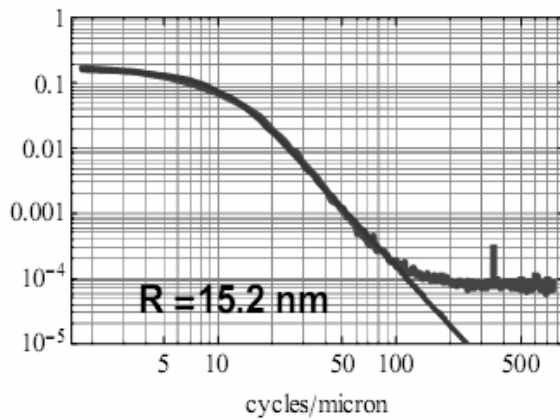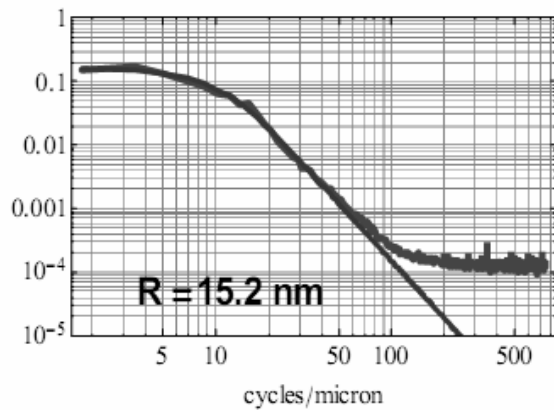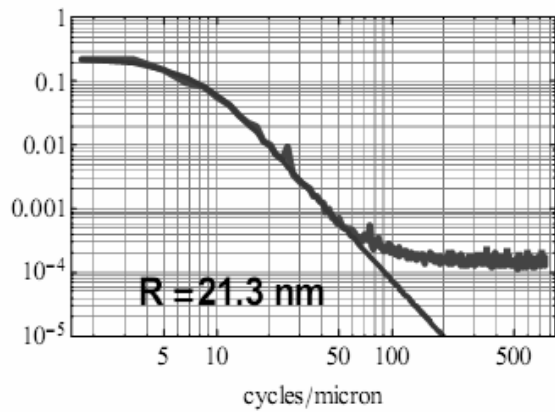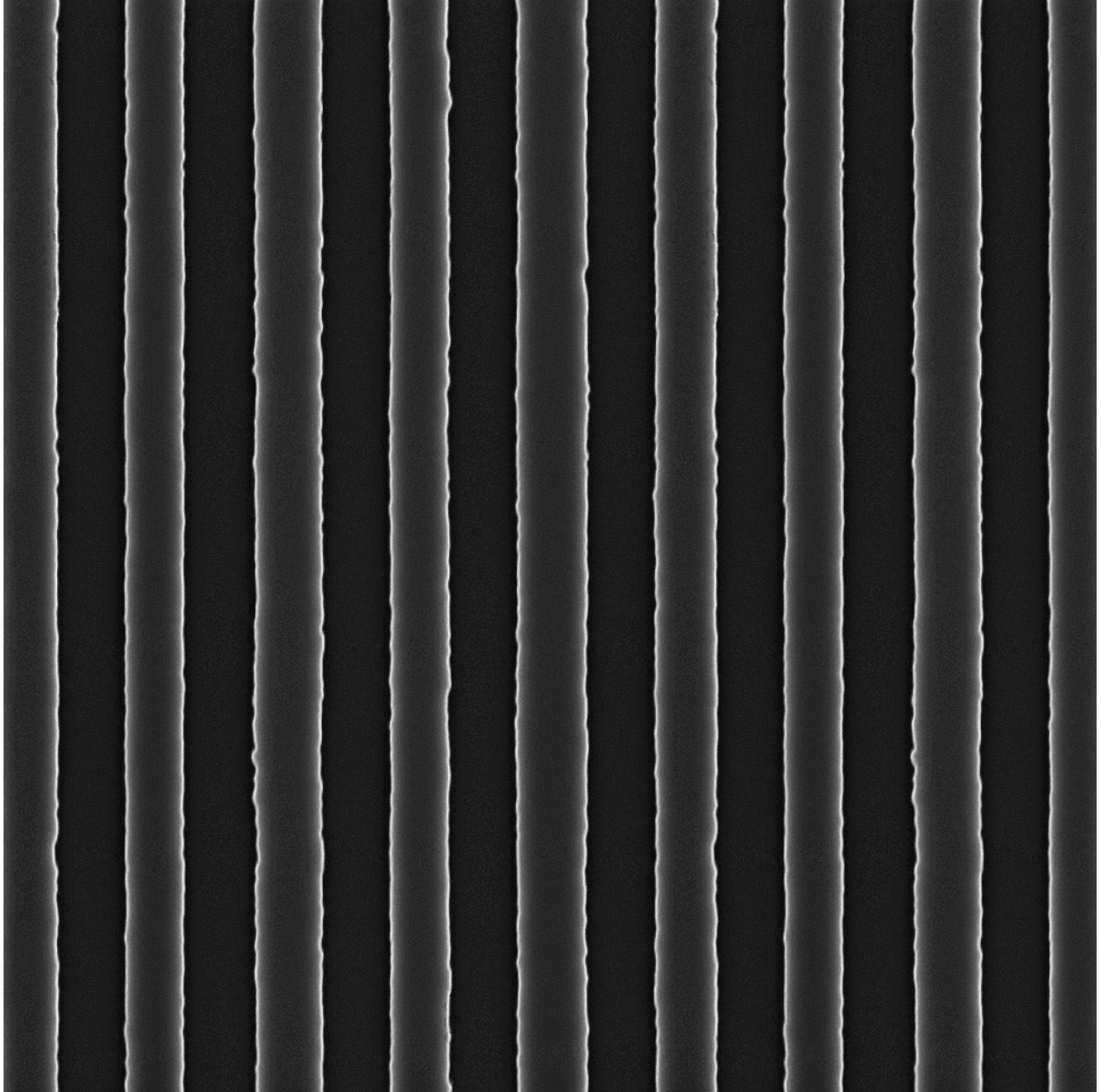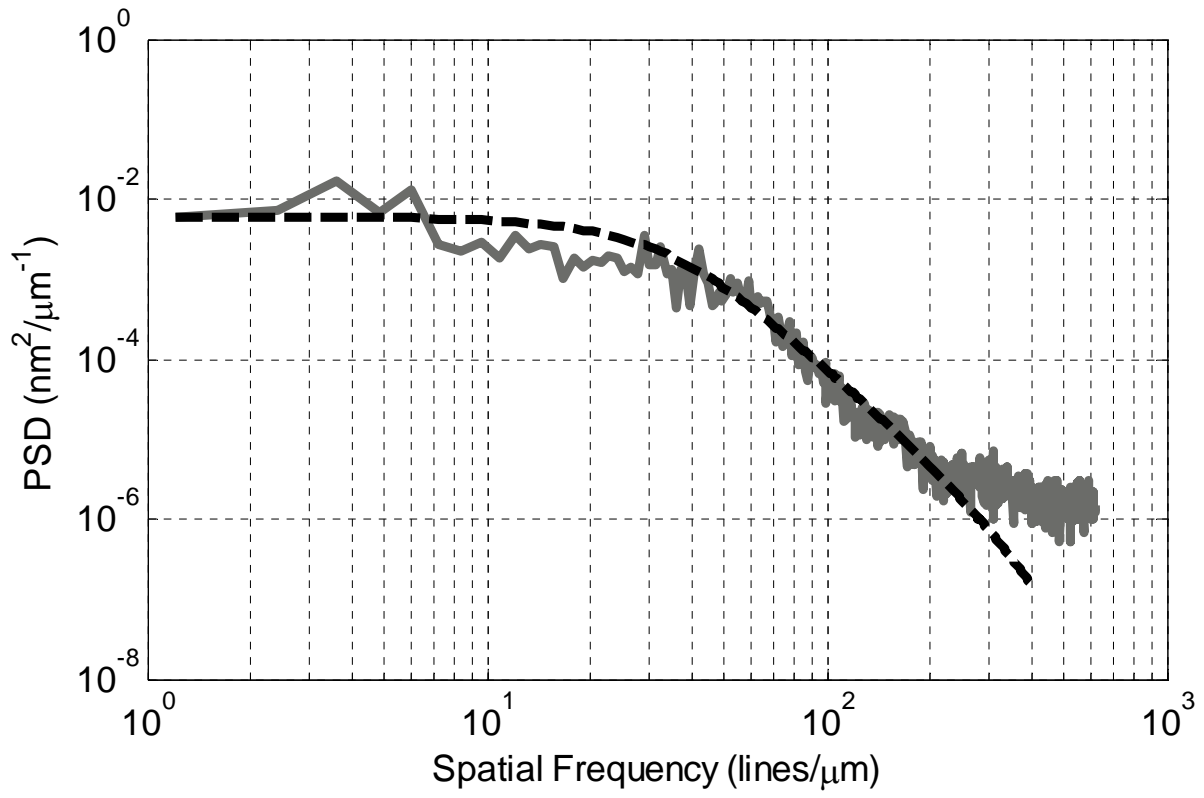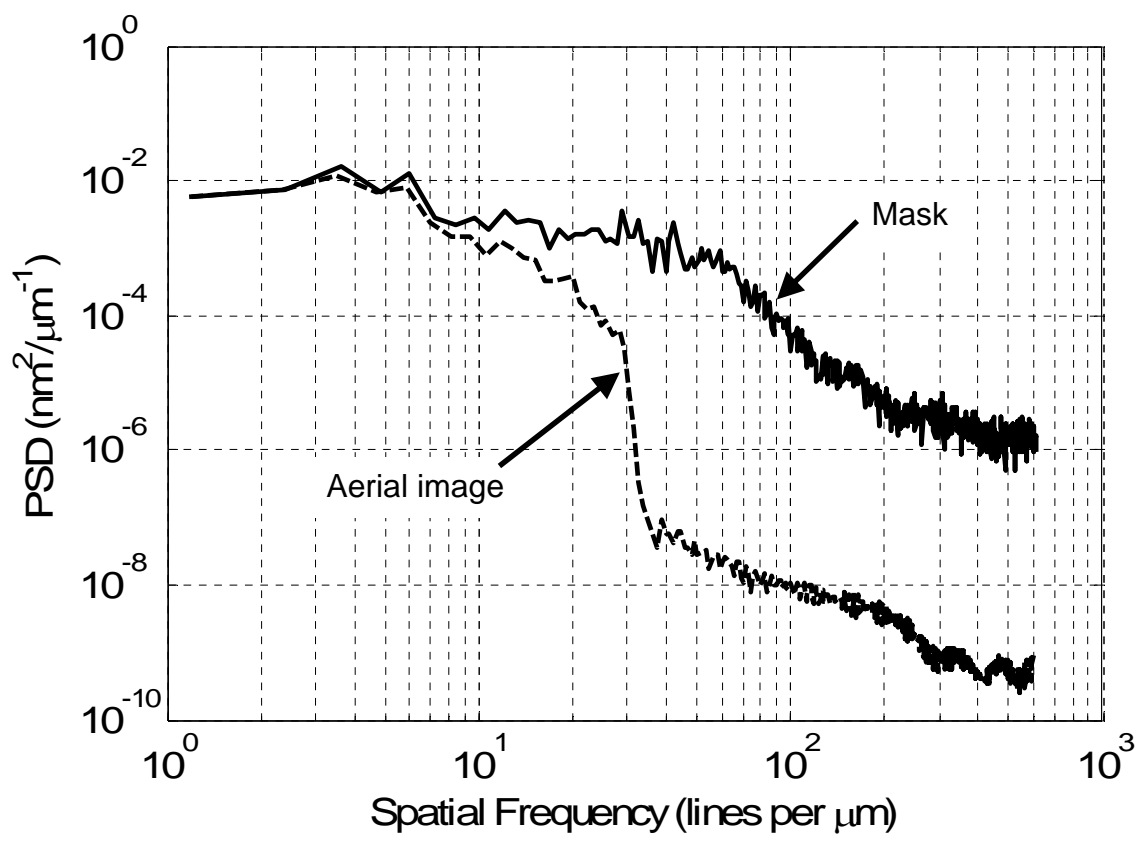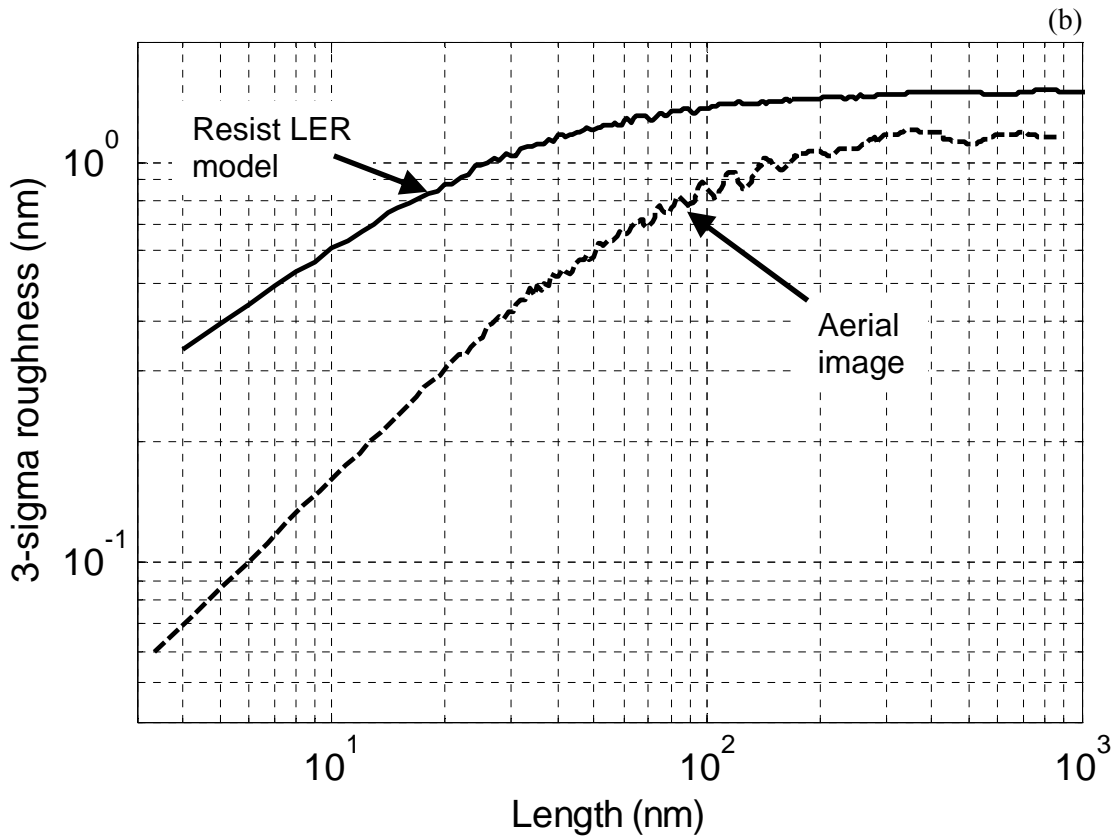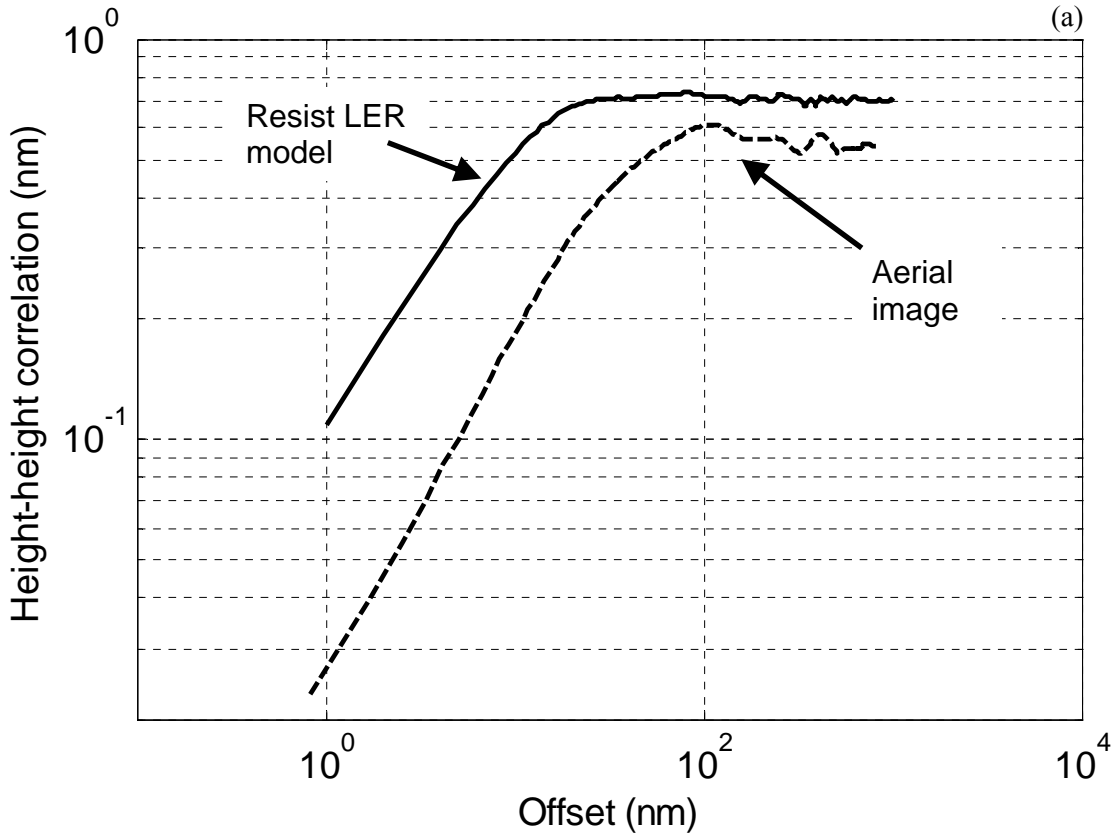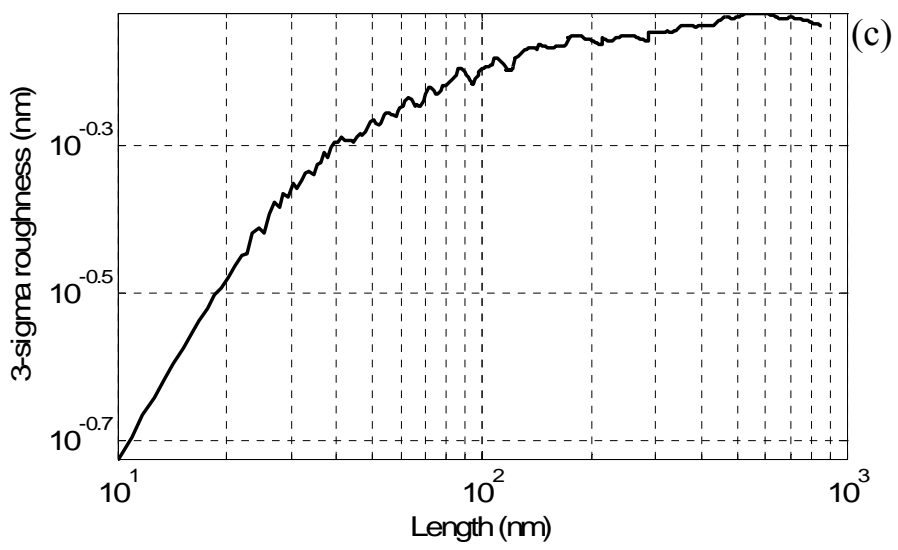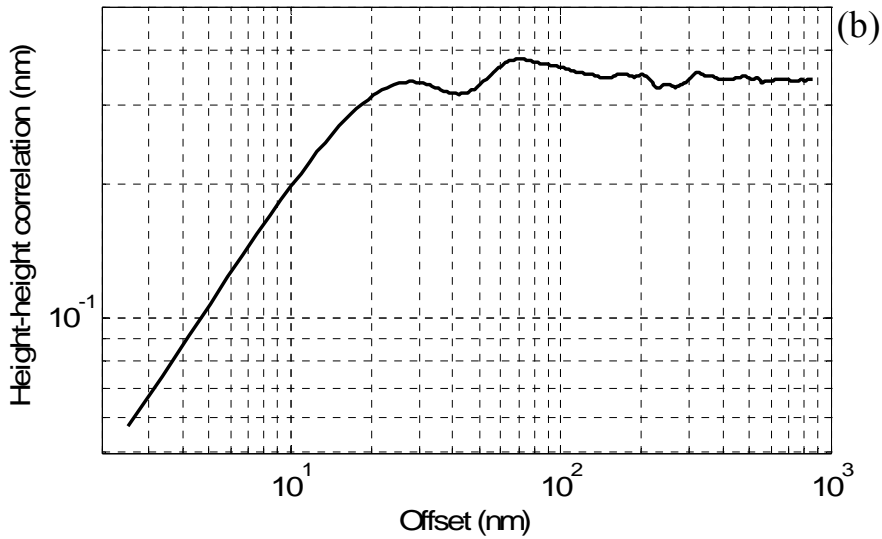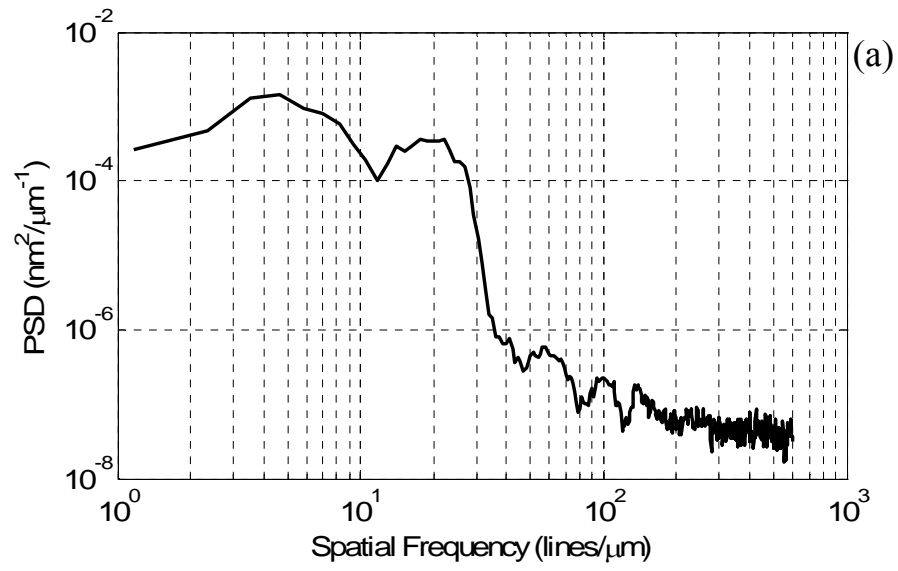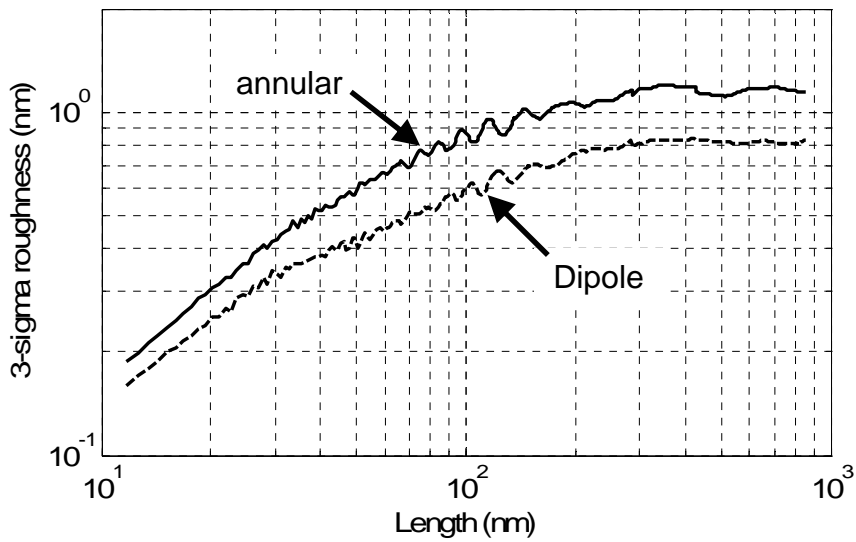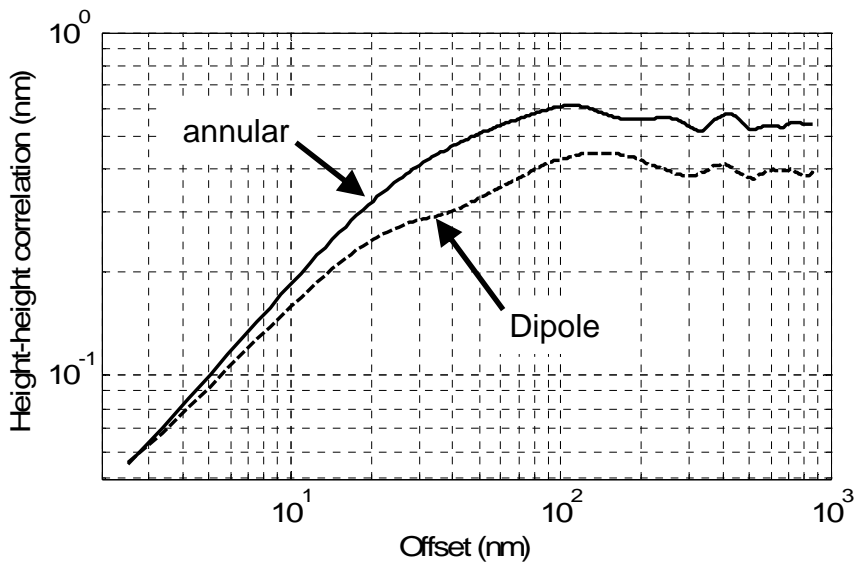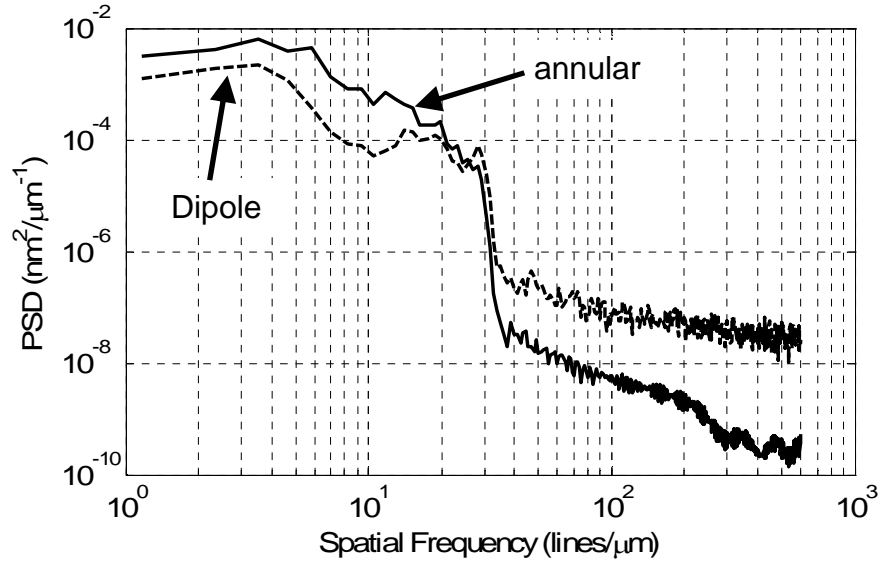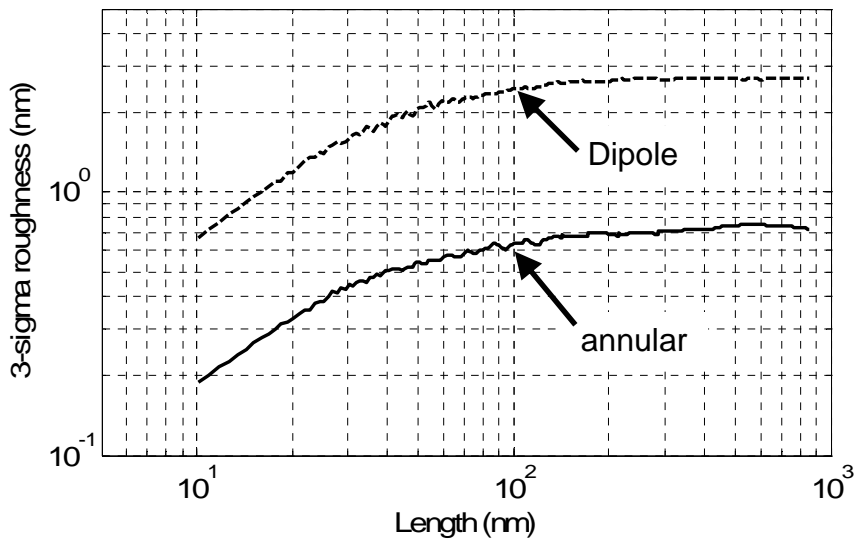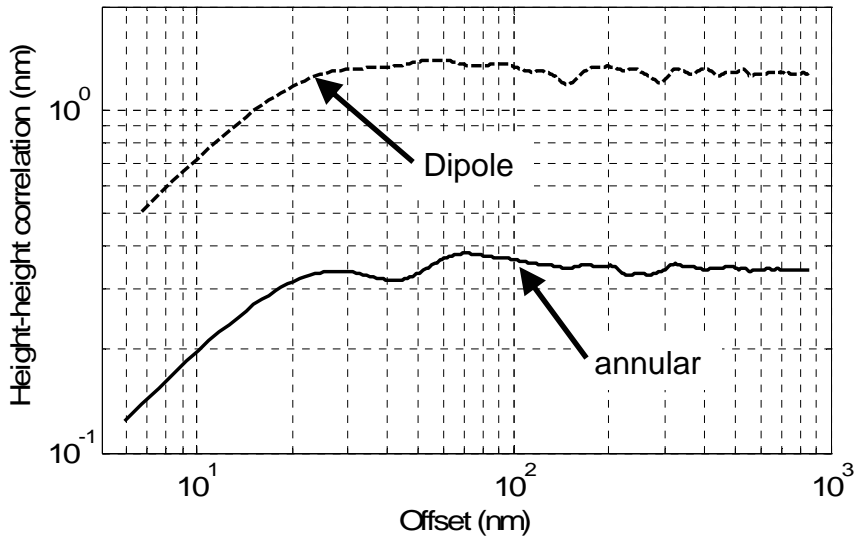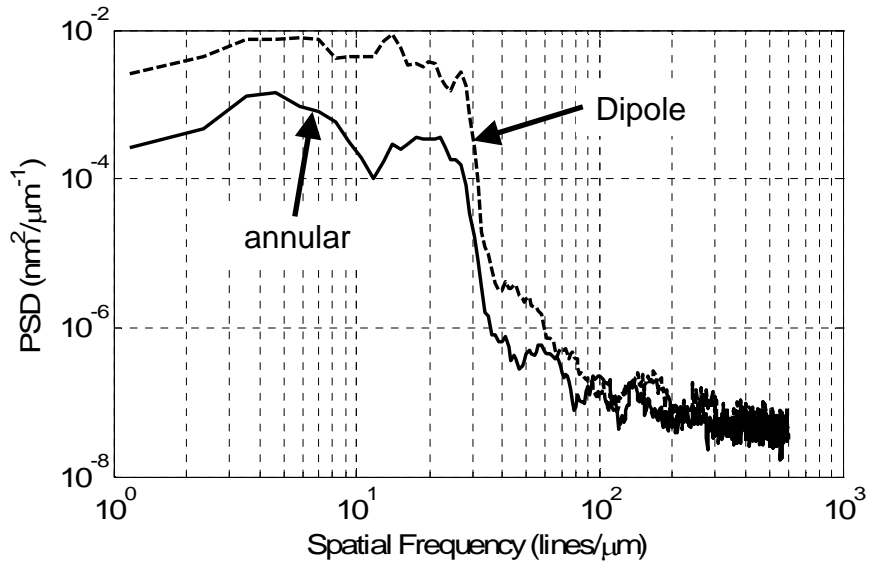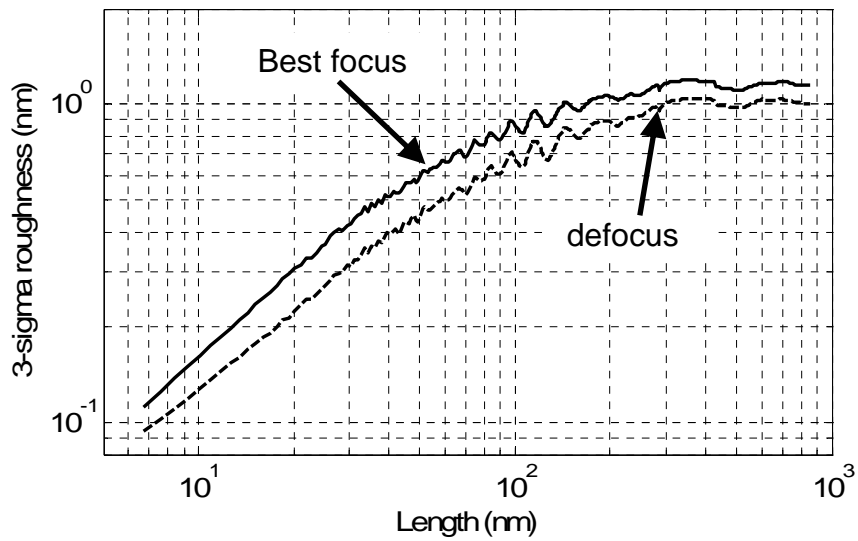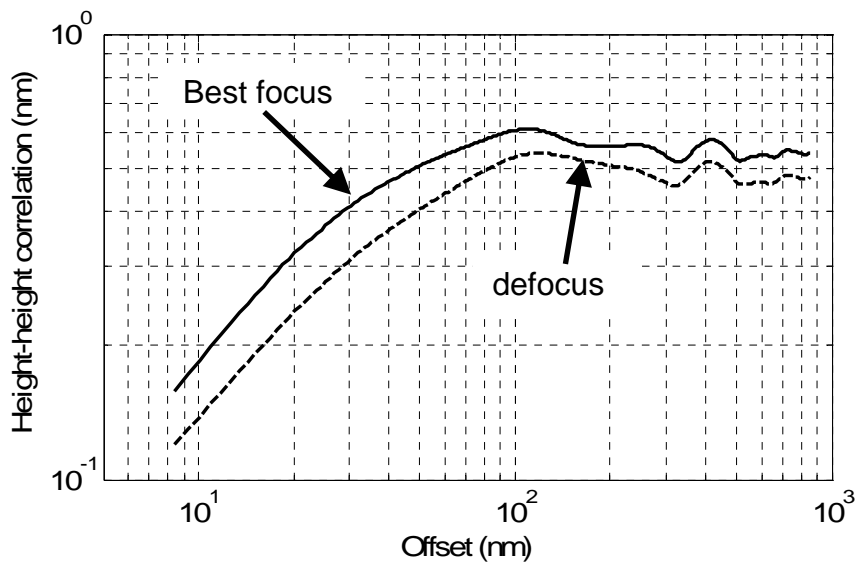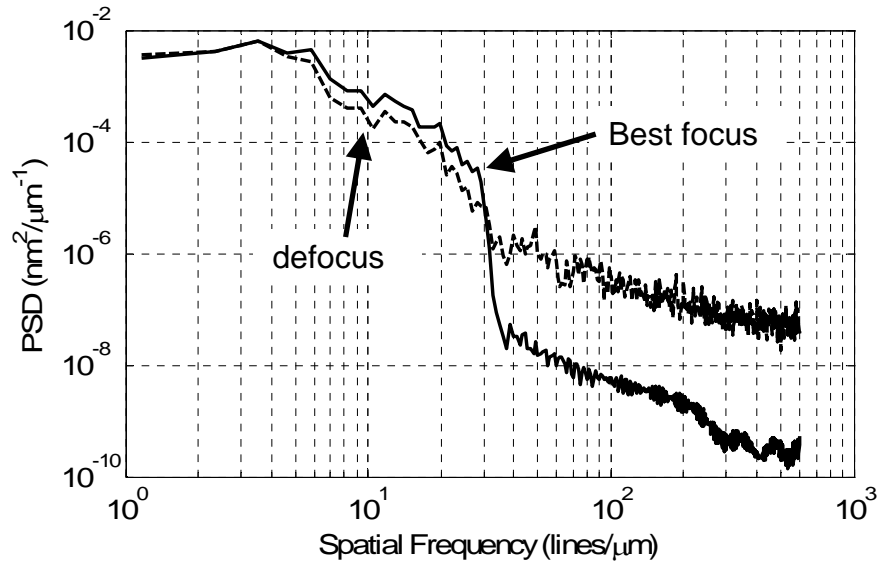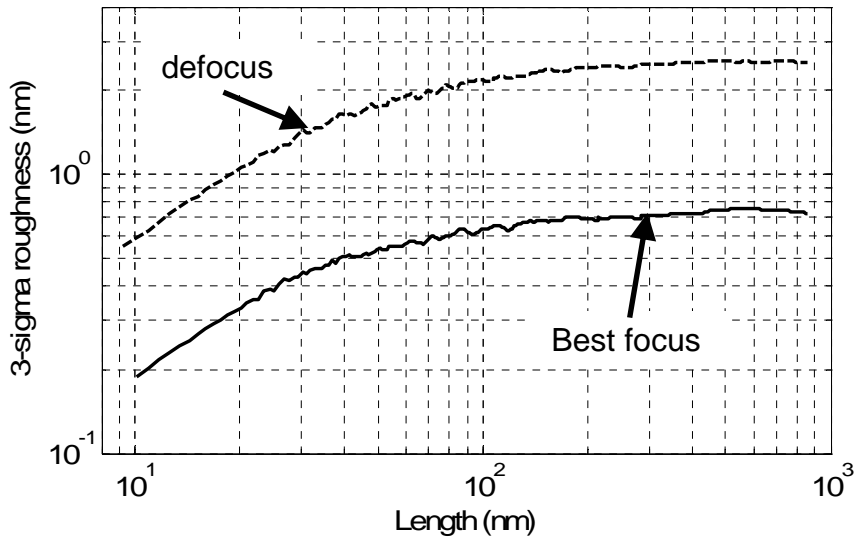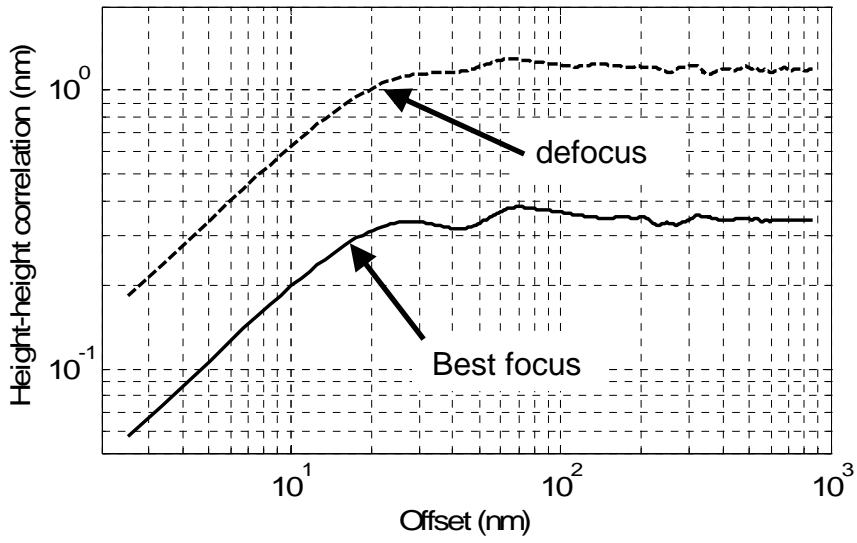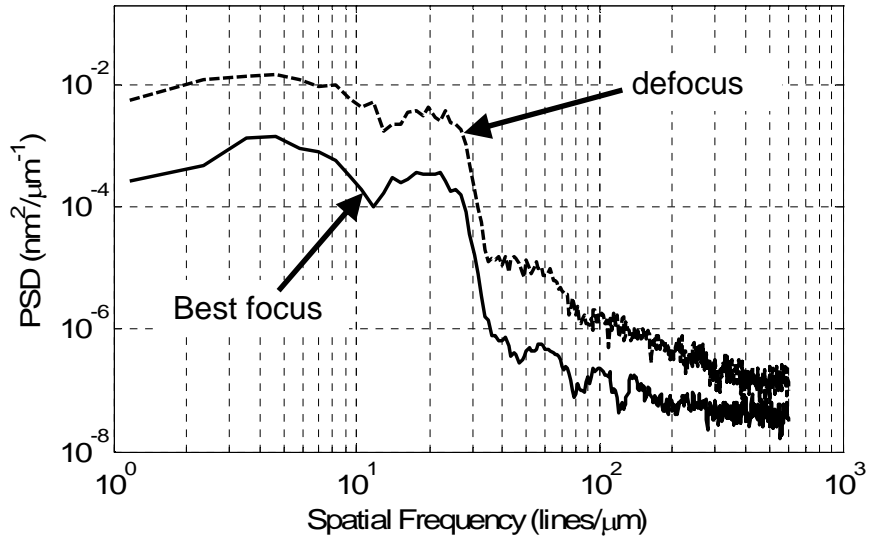
Naulleau, et al., Table 1

	Mask LER				Mask Multilayer Roughness			
	HHCF		PSD		HHCF		PSD	
	Lc (nm)	α	Lc (nm)	α	Lc (nm)	α	Lc (nm)	α
Annular	34.5	0.85	33.8	5	16.2	0.97	33.2	6
Dipole	38.4	0.82	31.6	5	16.2	0.91	32.1	6

Naulleau, et al., Table 2

	Mask LER				Mask Multilayer Roughness			
	HHCF		PSD		HHCF		PSD	
	Lc (nm)	α	Lc (nm)	α	Lc (nm)	α	Lc (nm)	α
Best Focus	34.5	0.85	33.8	5	16.2	0.97	33.2	6
110-nm defocus	47.8	0.83	31.6	1.9	18.8	0.88	32.8	3.2