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A Superconducting Focusing Solenoid for the Neutrino Factory Linear Accelerator*

M. A. Green, V. Lebedev, and B. P Strauss

Abstract--The proposed linear Accelerator that accelerates muons from 190 MeV to 2.45 GeV will use superconducting solenoids for focusing the muon beam. The accelerator will use superconducting RF cavities. These cavities are very sensitive to stay magnetic field from the focusing magnets. Superconducting solenoids can produce large stray fields. This report describes the 201.25 MHz acceleration system for the neutrino factory. This report also describes a focusing solenoid that delivers almost no stray field to a neighboring superconducting RF cavity.

Index Terms—Superconducting Focusing Solenoid, Active Shielding, Superconducting RF Cavities

I. INTRODUCTION

The proposed neutrino factory generates neutrinos from muon decay in a muon storage ring[1,2]. The muons are created from the decay of pions produced from the collision of protons on a fixed target. The pions are captured by a 20 T solenoid. The pions then decay to muons as they pass down a long channel. The muons are phase-rotated, bunched, and cooled before they can be accelerated and stored in the storage ring[3].

The first stage of acceleration is a linear accelerator that uses superconducting RF cavities[4]. Muon focusing during this acceleration occurs within superconducting solenoids. After the muons are accelerated to 2350 MeV, they enter a recirculating linac that accelerates them to their final energy of 20 GeV. The 20 GeV muons are then stored in storage ring where they decay to two neutrinos and an electron or positron (depending on the charge of the originating muon type).

II. FOCUSING SOLENOIDS FOR THE LINEAR ACCELERATOR

The proposed linear accelerator that accelerates the muons from 190 MeV to 2350 MeV extends nearly 400 meters. The Linear accelerator will use high gradient superconducting RF cavities to accelerate the beam. These cavities will not only increase the acceleration gradient, but they will also save electric power. Between the RF cavities there will be superconducting solenoids that provide beam focusing.

The 400-m long linear accelerator is divided into three sections. Upstream from the linear accelerator is a matching section that provides the first 60 MeV of muon acceleration. The end of the matching section consists of a pair of low stray field solenoids with adjustable currents for tuning. The first section has 4 short modules that are 5-meters long. Each short module has two cells of RF cavities. The second section has 16 medium length modules that are 8-meters long. Each intermediate module has two cells of RF cavities arranged in groups of two with a 1-meter space between groups. The short and intermediate modules have a 1-meter long focusing solenoid with a 1.25-meter long straight section at each end. The third acceleration section has 19 long modules that are 13-meters long. The long module has eight superconducting RF cells arranged in groups of two with a 1-meter space between cell groups. The focusing solenoid is 1.5 meters long with a 1.25-meter long space separating it from the adjacent RF cavity cells. The three types of acceleration modules used for the linear accelerator are illustrated in Figure 1.

Unlike focusing quadrupoles, focusing solenoids can produce a stray field that exists some distance from the magnet. The superconducting RF cavities are sensitive to magnetic fields even at the level of 0.0001 T (versus 0.043 T for the smallest single solenoid in Table 1), so a key parameter in the focusing solenoid design is getting rid of the solenoid stray field in the RF cavities. The methods one can use to eliminate the stray field in the RF cavities are:1) The focusing solenoid should produce zero net magnetic moment[5,6]. This means that the coil that produces the field is bucked by a coil that is larger in diameter. 2) The field from the bucking coils should be distributed in the same way as the solenoid field. This suggests that the bucking solenoid be around the focusing solenoid so that the return flux from the focusing solenoid is returned between the focusing solenoid and the bucking solenoid[7]. 3) The nested solenoids should be surrounded by iron except where the muon beam passes through the solenoid. 4) An iron flux shield should be installed between the focusing solenoid and the RF cavities. 5) If the stray field is still not low enough, the RF cavities can be covered with a type 2 superconducting shield. This shield will trap earth's magnetic field, but it will shield out the remaining stray flux from a nearby solenoid provided the cavities are cold (below 5 K) when the solenoid is turned on. An iron shielded actively shielded solenoid appears to be a minimum requirement in order to reduce the stray field to an acceptable level.

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Long Acceleration Cell (13-m Long)

Fig. 1. The Three Linear Acceleration Cells for the Superconducting Linac. Shown in the figure are the superconducting RF cavities and the supperconducting solenoid magnets used for focusing. No detail is shown concerning the solenoids.

TABLE 1.	PARAMETERS FOR	THE FOCUSING	SOLENOIDS
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	Short Cell	Medium Cell	Long Cell
Beam Bore Diameter (mm)	460	460	300
Iron Shell Length (mm)	1300	1300	1750
Iron Shell Outer Diameter (mm)	1240	1240	1120
Iron Shell Thickness (mm)	9.5	9.5	9.5
Inner Coil Average Radius (mm)	254	254	182
Inner Coil Thickness (mm)	10.4	10.4	31.2
Number of Inner Coil Layers/Turns	8/4840	8/4840	24/21816
Outer Coil Average Radius (mm)	520.6	520.6	453.6
Outer Coil Thickness (mm)	2.6	2.6	5.2
Outer Coil Center Gap (mm)	50	50	50
Number of Outer Coil Layers/Turns	2/1152	2/1152	4/3512
Solenoid Cold Mass (kg)	376	376	746
Solenoid Cryostat Mass (kg)	166	166	238
Iron Shell Mass (kg)	485	485	581
Solenoid Average Magnetic Induction (T)	2.1	2.1	4.2
Solenoid Magnetic Length (m)	~1.0	~1.0	~1.5
Magnet Design Current (A)	469.6	469.6	274.0
Peak Induction in the Inner Coil Bp (T)	~2.9	~2.9	~5.8
Magnet Conductor Ic at 4.4 K and Bp (A)	~1100	~1100	~590
S/C Current Density J (MA m ⁻²)	307	307	180
Solenoid Stored Energy E (MJ)	0.421	0.421	1.306
Solenoid Self Inductance (H)	3.82	3.82	34.8
Quench Protection E J^2 Limit ($A^2 m^4 J$)	3.97×10^{22}	3.97×10^{22}	4.23×10^{22}

Fig. 1 shows the location of iron flux barriers that provide additional magnetic shielding for the nearby RF cavities. The barrier is designed to remove most of the magnetic flux that passes through the beam hole in the iron shield around the superconducting focusing solenoid.

Table 1 presents the mechanical and electrical parameters for the short, medium, and long module focusing solenoids. The matching solenoids at the start of the channel are similar to the first focusing solenoids. All of these solenoids are designed to have zero net magnetic moment. The solenoids in Table 1 are assumed to have a warm bore and a warm iron shell around the solenoid pair. It should be noted that the magnet bore does not have to be warm. The solenoid pair is assumed to be powered through a single set of HTS and gascooled electrical leads. Since the magnet pair is hooked in series, the focusing solenoids will have nearly zero net magnetic moment at all excitation currents.

The design of the solenoid and its bucking coil is a case where the cost of the magnet is to be minimized. A bucking coil that has too small a diameter will create a large bucking field. More superconductor is needed in the inner solenoid in order to produce the desired on axis field. A bucking coil that is too large will result in a magnet system with more overall mass even though that magnet may have a smaller amount of superconductor within it. A case that is close to optimum is a case where the bucking coil has twice the radius of the inner coil. The magnets shown in Table 1 were designed to have an even number of layers of superconductor in both coils. There is a gap in the outer coil that allows one to attach current leads.

Fig. 2 shows a half cross section of the long focusing solenoid (1.5-m long with 4.2 T in the inner bore) in a plane that goes along the solenoid axis. Fig. 2 shows the separation of the inner coil and the bucking coil. Also shown are the magnet cryostat, an electrical lead, and the iron shield around the actively shielded solenoid. The center of the cryostat has no iron shield around it because there is very little magnetic flux leaking outside the bucking solenoid. The magnet shown in Fig. 2 is two concentric solenoids. A superconducting pancake at the end of the magnet may be more effective in shielding the nearby RF cavity. In any event the net magnetic moment should be zero.

The solenoid shown in Fig. 2 has the inner solenoid mounted on the inside of an aluminum shell. The outer coil is wound on the outside of an aluminum bobbin.

The conductor proposed for the focusing solenoid insulated dimensions of 1.0 by 1.65 mm. This is a standard MRI conductor with a copper to superconductor ratio of 4. The conductor contains fifty-five 80-micron filaments of niobiumtitanium. The conductor twist pitch is 12.7 mm. The proposed conductor has a guaranteed critical current of 760 A at 4.2 K and 5 T. The turn to turn insulation on the proposed conductor is a varnish. The layer to layer insulation proposed for the magnet is about 0.35 mm of epoxy impregnated fiberglass. The magnet coil can either be wound and then vacuum impregnated or it could be wound as a wet lay-up with a room temperature cure.



Fig. 2. A Cross Section of the 1.5-meter Long Focusing Solenoid Parallel to the Solenoid Axis

III. COOLING OF THE FOCUSING SOLENOID

The focusing solenoids are cooled by conduction from the 6061-aluminum support structure. The aluminum support structure will be cooled by two-phase helium flowing in tubes attached to it. It is assumed that all of the magnets in an acceleration section are cooled in series from the two-phase helium refrigerator and control cryostat. Whether this refrigerator is the same one that cools the superconducting RF cavities depends on the operating temperature of the RF cavities. The 4.4 K heat load is about 0.5 W for the short solenoid. The long solenoid heat leak is about 0.7 W. One can cool 20 magnets with a two-phase helium a mass flow rate of 2.5 g s⁻¹. The two-phase helium cooling tubes will be attached to the coil support structures, the attachment points for the cold mass supports, and the base of the HTS leads.

The heat load into the shield circuit helium stream is expected to vary from 6.1 to 7.3 W depending on the length of the magnet. The shield gas comes from the refrigerator at a temperature of 30 K. This gas enters the magnet cryostat through a single vacuum insulated tube. The helium flow in this tube is dictated by the needs of the gas-cooled leads between 50 K and room temperature. The mass flow through the shield circuit is governed by the needs of the gas-cooled leads. The short solenoid leads will need 0.05 grams per second; the long solenoid leads will need 0.035 grams per second. The gas exits from the gas-cooled electrical leads at room temperature. It returns warm to the refrigerator compressor suction. In the short solenoid, the shield gas enters the gas-cooled leads at about 55 K. In the long solenoid, the top of the HTS leads will be about 70 K. The same HTS leads can be used for both magnets.

IV. POWER SUPPLY AND QUENCH PROTECTION

For simplicity sake, it is assumed that each focusing solenoid has its own power supply. A 5 V and 500 A power supply can be used for both magnets. The charge time for the short magnet, with 3 V across its leads is about 600 second. The charge time for the long magnet is about 3200 seconds with 3 V across its leads. The power supply should allow one to charge and discharge the magnet and control its current to a set value. The magnet approver supply is not required to operate the magnet at both positive and negative currents.

Quench protection for the magnet is by quench-back from the support structure and bobbin. The proposed quench protection system is completely passive. As a normal region develops in one coil, a current is induced in the adjacent coil bobbin. This current heats the superconducting coils driving them normal. Since each magnet is individually powered, a quench back quench protection system is more than adequate to protect the focusing solenoids.

V. CONCLUDING COMMENTS

The focusing solenoids for the neutrino factory linear accelerator can be built so that stray magnetic flux does not affect the superconducting RF cavities that are about 1-meter from the focusing solenoid. The superconducting solenoid can be built so that it generates nearly zero net magnetic moment at all excitations. This is the principle behind the actively shielded solenoids commonly used for MRI magnets and NMR magnets. Active shielding (zero net magnetic moment) alone will not keep stray field from the magnet from affecting the RF cavities. The distribution of the currents in the bucking coil is also important. An iron shield around the focusing solenoid will remove all of the stray flux from the magnet except for flux, which passes through the beam hole. Iron shields that are about 300 mm from the end of the focusing solenoid iron shield will further reduce the magnetic flux seen by the RF cavity.

The superconducting focusing solenoids can be built using commercial MRI magnet conductor. The superconducting coils will be wound and potted. They will be indirectly cooled by conduction from the support structure. The support structure will be cooled using two-phase helium carried in tubes that are attached to the support structure or bobbin.

As with modern MRI magnets, the focusing solenoids would have HTS leads between the magnet at 4.4 K and the shield at say 50 K. Helium gas at 30 K is used to cool the magnet shield. After the shield gas cools the shield and cold mass support intercepts it can be used to cool the gas-cooled current leads that are attached to the top of the HTS current leads.

Each focusing solenoid will be separately powered. The actively shielded solenoids are protected during a quench by quench-back from the support structure and bobbin.

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