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Tomography of Southern California via Bayesian Joint Inversion of Rayleigh Wave Ellipticity and Phase Velocity from Ambient Noise Cross-Correlations

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Key Points:

- Ambient noise cross-correlations across Southern California show clear Rayleigh waves with measurable phase and amplitude information
- Amplitude information, through Rayleigh ellipticity (H/V), gives new constraint on nearsurface structure due to shallow crust sensitivity
- Shallow crust structures previously only seen in local studies are recovered regionally via joint inversion of phase velocity and H/V

1 Abstract

A self-consistent regional-scale seismic velocity model with resolution from seismogenic depth 2 to the surface is crucial for seismic hazard assessment. Though Southern California is the most 3 seismically imaged region in the world, techniques with high near-surface sensitivity have been 4 applied only in disparate local areas and have not been incorporated into a unified model with 5 6 deeper resolution. In the present work, we obtain isotropic values for Rayleigh wave phase 7 velocity and ellipticity in Southern California by cross-correlating daily time-series from the year 2015 across 315 regional stations in period ranges 6 to 18 seconds. Leveraging the 8 complementary sensitivity of the two Rayleigh wave datasets, we combine H/V and phase 9 velocity measurements to determine a new 3D shear velocity model in a Bayesian joint inversion 10 framework. The new model has greatly improved shallow resolution compared to the SCEC 11 CVMS4.26 reference model. Well-known large-scale features common to previous studies are 12 resolved, including velocity contrasts across the San Andreas, San Jacinto, Garlock, and Elsinore 13 faults, mid-crustal high-velocity structure beneath the Mojave Desert, and shallow Moho beneath 14 the Salton Trough. Other prominent features that have previously only been imaged in focused 15 local studies include the correct sedimentary thickness of the southern Central Valley, fold 16 structure of the Ventura and Oak Ridge Anticlines, and velocity contrast across the Newport-17 Inglewood fault. The new shallow structure will greatly impact simulation-based studies of 18

19 seismic hazard, especially in the near-surface low-velocity zones beneath densely populated

20 areas like the Los Angeles, San Bernardino, and Ventura Basins.

21 **1 Introduction**

Southern California is one of the most tomographically imaged regions in the world from 22 a variety of methods including local body waves (e.g., Lin et al., 2010; Allam & Ben-Zion, 23 2012), teleseismic body waves (Schmandt & Humphreys, 2010), surface waves (Tanimoto & 24 Prindle Sheldrake, 2002; Yang & Forsyth, 2006), ambient noise (e.g., Lin et al., 2008; Zigone et 25 al., 2015; Barak et al., 2015), and full waveforms (e.g., Chen et al., 2007; Tape et al., 2009). The 26 most recent models combining earthquake and ambient noise data (Lee et al., 2014; Fang et al., 27 2016) are extremely detailed and can successfully replicate observed earthquake waveforms at 28 relatively high frequency (Taborda et al., 2016). However, one of the main limitations of these 29 models is that they provide only weak constraint on the uppermost crustal structure (<3km) due 30 to the relatively long periods employed (above a few seconds) and because amplitude 31 information is not included. Shallow structure is well constrained by recent, focused active-32 source seismic or local earthquake double-difference tomography studies in a few sub-regions 33 (e.g., Fuis et al., 2001; Süss & Shaw, 2003; Allam et al., 2014; Fuis et al., 2017), but is lacking 34 regionally. Improved models of upper crustal structure are crucial because they allow vastly 35 improved seismic ground motion predictions (Vidale and Helmberger, 1988; Graves et al., 36 2011), ameliorate misinterpretations of mantle structure (Waldhauser et al., 2002; Bozdağ and 37 Trampert, 2008; Schulte-Pelkum & Ben-Zion, 2012), provide insight into lithospheric 38 discontinuities (Langston, 2011), and validate geological interpretations based on surface 39 observations (e.g., Graymer et al., 2005). Because Rayleigh wave horizontal-to-vertical 40 (hereafter, denoted as H/V for conciseness) ratios have shallower sensitivity than phase velocity, 41 they can provide much stronger constraints on shallow crustal structure at regional scales 42 (Tanimoto & Rivera, 2008; Lin et al., 2012; Lin et al., 2014). 43

Surface wave tomography using phase or group velocities measured from ambient 44 seismic noise cross-correlations is by now a standard technique (e.g., Bensen et al., 2007; Lin et 45 al., 2009; Campillo et al., 2011). Traditional methods for measurement of noise spectral H/V 46 ratio (e.g., Nakamura, 1989; Fäh et al., 2001) have been used to image structure in the upper few 47 hundred meters, characterize site response, and predict ground motion. Interpretation of these 48 measurements, however, depends on the assumed noise character, i.e. Rayleigh wave dominant, 49 body wave dominant, or a mix (Bonnefoy-Claudet et al. 2006). Rayleigh waves isolated from the 50 noise wavefield using noise cross-correlation, on the other hand, have been used to measure short 51 period Rayleigh wave H/V ratios and recover shallow velocity structure across the U.S. (Lin et 52 al., 2014). Recently, Rayleigh H/V ratios have been measured in Southern California and used to 53 interpret shallow structure (Muir and Tsai, 2017). However, a joint 3D inversion for regional 54 shear velocities incorporating both H/V data to constrain shallow structure and phase velocities 55 to constrain the mid-crust has not been conducted. 56

57 In the present work, we leverage this complementary sensitivity of the Rayleigh wave amplitude and phase data to image the crustal shear wave velocity structure throughout Southern 58 California. We perform detailed analysis of Rayleigh wave H/V ratio from noise cross-59 correlations on 315 stations and combine H/V and phase velocity measurements in a Markov 60 Chain Monte Carlo joint inversion. This method benefits from the quantification of full model 61 62 uncertainties and analysis of misfit while avoiding local minima by testing an ensemble of candidate models (Shen et al., 2012; Shen & Ritzwoller, 2016; Roy & Romanowicz, 2017). In 63 Section 2, we describe the cross-correlation and measurement of H/V ratios, quality control 64 criteria, and the Markov Chain Monte Carlo inversion methods. We present the results in Section 65 3 and discuss in detail the relevance of our final shear wave velocity model to the complicated 66 geology of Southern California in Section 4. 67

68

69 2 Data & Methods

70 2.1 Data and Stations

71 From the Southern California Earthquake Data Center we obtain a year of continuous waveforms from 315 three-component stations available for 2015 in the Southern California 72 plate boundary region (Figure 1). These stations are associated with multiple seismic networks 73 including the Anza network, Southern California Seismic Network, and the San Jacinto Fault 74 Zone network. From these networks, we incorporate multiple seismic instruments including 75 broadband (BH, HH, BN, HN) and short period (EH). This data set allows us to analyze several 76 77 geological regions of interest including the Coast Ranges, Central Valley, Mojave Desert, Sierra Nevada Range, Los Angeles Basin, San Jacinto Fault Zone, Peninsular Ranges, and the Salton 78 Trough (Figure 1). Our study area also includes the Transverse Ranges, between the Coast 79 Ranges, LA basin and the Mojave Desert, although this is not marked in Figure 1 in order to 80 emphasize station coverage. 81

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- 83

2.2 Ambient Noise Pre-Processing and Cross-Correlation

We closely follow the method described in Lin et al. (2014) to process the daily noise time series for each station prior to cross-correlation and stacking. For each station, data is cut into daily noise time series followed by decimation to a sampling rate of 4 Hz. We remove the mean, trend and instrument response for each east, north, and vertical component (E, N, Z) and apply a bandpass filter between 5 s and 150 s periods. Following Bensen et al. (2007), we next

remove earthquake signals and instrumental irregularities via temporal normalization. To obtain 89 90 the temporal normalization functions, we bandpass the seismic signal between 15 and 50 second period and calculate a 128-s time window running absolute mean for each component. For each 91 92 point in time, we divide the three-component unfiltered time-series by the maximum of all temporal normalization functions for that corresponding time across all components (E, N, Z). 93 Following temporal normalization, we perform spectral whitening to broaden the period band to 94 increase potential recovery of surface-wave signals (Lin et al., 2008), dividing the spectrum of 95 each component by the average amplitude of the three-component (E, N, Z) smoothed, 0.025 Hz 96 (or 20 points halfwidth) running-mean, spectra. By applying the same temporal normalization 97 and spectral whitening to each component (E, N, Z) we preserve the relative amplitude between 98 components and allow rotation to be applied after cross-correlation and stacking. Due to the 99 commutative nature of these normalization processes, we perform these steps pre-cross-100 correlation, which saves significant computational cost (Lin et al., 2008). 101

We next apply the methods of Lin et al. (2008) to calculate the nine-component cross-102 correlations among the north, east and vertical components, stack all daily cross-correlations 103 from 2015, and rotate the horizontal motion into radial (R) and transverse (T) directions. After 104 105 rotation, we analyze the positive time lag (causal) and negative time lag (acausal) parts independently to avoid potential mixing of good and bad signals. Figure 2a shows an example of 106 the ZZ, ZR, RZ and RR cross-correlations bandpassed around 8 s period between a station in the 107 108 LA Basin (FMP) and a station in the Sierra Nevada (WBS), highlighted in yellow and red, respectively, in Figure 1. Clear Rayleigh waves with similar arrival times are seen on all four 109 components but the observed amplitudes are different, indicating contrasting horizontal-to-110 vertical ratios at each station site. The impact of asymmetric distribution of noise sources, normal 111 to the coastline, due to oceanic waves dominating the noise field (Hillers et al., 2013) creates 112 stronger signal on the positive time lag in Figure 2a due to the location of the virtual source 113 (FMP) in relation to the receiver (WBS). We are still able to retrieve the empirical Green's 114 function, detected through signal-to-noise ratio, with inhomogeneous noise distribution due to 115 sufficiently strong ambient noise and its natural scattering properties (Yang & Ritzwoller, 2008; 116 Lin et al., 2008). 117

The negative time derivative of the ZZ, ZR, RZ and RR cross-correlations, assuming a 118 diffuse noise wavefield (Lobkis & Weaver, 2001), is related to the Rayleigh-wave Green's 119 functions for a point force in the vertical (Z) or radial (R) direction at the source station and 120 121 recorded in the vertical and radial directions at the receiver station. As shown in Figure 2b, by combining ZZ and ZR cross-correlations or RZ and RR cross-correlations, we can study the 122 Rayleigh-wave particle motion at the second (acting as the receiver) station. Similarly, by 123 combining ZZ and RZ cross-correlations or ZR and RR and considering the reciprocity of the 124 Green's function (Aki & Richards 2002), we can analyze Rayleigh-wave particle motion at the 125 first (acting as the source) station. As the wave is traveling from the source, FMP, to receiver, 126 127 WBS, the receiver station shows retrograde particle motion. In contrast, the particle motion is prograde for the particle motion associated with the source station. 128

From Figure 2b it is apparent that the amplitude ratios differ strongly between the two stations; receiver station WBS is elongated vertically (low H/V), and receiver station FMP is elongated horizontally (high H/V). This is due to the highly localized sensitivity of Rayleigh wave horizontal to vertical (H/V) amplitude ratios to shallow structure. Station FMP is located in the LA Basin, where the contrast of soft sediment and bedrock elongates the radial component and creates a high horizontal to vertical (H/V) amplitude ratio. Receiver station WBS is located on crystalline rock, which without a pronounced shallow to deep velocity contrast creates a low

136 H/V ratio. In order to image both shallow and mid-crustal structure throughout the area, we make

137 measurements of both H/V and phase velocities from the cross-correlation functions over period

ranges of 6-18 s and 6-16 s, respectively. The relative depth sensitivities for shear waves of

139 phase velocities and H/V ratios are shown by their sensitivity kernels in Figure 3 for a location

140 on the San Andreas fault (Fig. 1 green star). These sensitivity kernels demonstrate the

- 141 complimentary sensitivity of the two Rayleigh wave measurements.
- 142

143 **2.3 H/V**

We use frequency-time analysis (FTAN; Bensen et al., 2007) to determine the maximum 144 amplitude of the envelope for both causal and acausal sides of the ZZ, ZR, RZ and RR cross-145 correlations. Next, we measure H/V independently on both the causal and acausal portions of the 146 correlograms. Specifically, for the first station (source station) we calculate H/V using RZ/ZZ 147 and RR/ZR cross-correlation amplitude ratios for both the causal and acausal signals (i.e. four 148 H/V measurements per station of each station pair). Similarly, for the second station (receiver 149 station) H/V is determined using ZR/ZZ and RR/RZ cross-correlation amplitude ratios. We only 150 retain good measurements by imposing several selection criteria, including signal-to-noise ratio 151 greater than 5 and interstation distance larger than three wavelengths to satisfy the far-field 152 condition (Bensen et al., 2007). For each period band, we define the signal-to-noise ratio as the 153 average of the ratio of peak energy within the expected Rayleigh wave signal window, between 154 1.5km/s and 4.5km/s, to the root mean square of noise before and after the expected signal 155 window (Lin et al., 2008). We apply this method to all cross-correlations to establish a large 156 number of H/V measurements for each station. 157

We further stabilize each station's result in a quality control process designed to remove spurious measurements. We iteratively remove all measurements greater than three standard deviations, recomputing the mean and number of measurements within two standard deviations until no further measurements are discarded. The final range of measurements after stabilization is marked for stations FMP and WBS in Figure 4 via dashed lines. For each station, we then use H/V measurements retained after stabilization to calculate isotropic H/V and uncertainty as the mean and standard deviation of the mean for each station, respectively.

Because H/V is a measurement of the relative change of two quantities, we express both 165 H/V and uncertainty as logarithms (Figure 5a,c); this has been shown to be the only symmetric, 166 additive, and normed indicator of relative change (Törnqvist et al., 1985). Since H/V is a local 167 168 measurement, we perform variable Gaussian smoothing to resolve log10(H/V) measurements throughout the entire region with 0.05°-by-0.05° spaced grid points (Fig. 5a,c) setting the 169 maximum Gaussian half-width as distance to the three nearest stations. To prevent overly-170 171 smoothed results, we discard any points within the region that do not have three stations within 172 50km. To propagate uncertainty, we determine the gaussian-weighted uncertainty from the standard deviation of the mean H/V of stations within the defined Gaussian distance for each 173 174 grid point (Fig. 5b,d).

175 176

2.4 Phase Velocities from Eikonal Tomography

We use the isotropic phase velocity (Vph) maps derived by Qiu et al. (2018) over the period range of 6 to 16 seconds. Rayleigh wave phase-velocity dispersion curves are first measured from vertical-vertical cross-correlations for all available station pairs in the same region as the present study (Figure 1). Eikonal tomography (Lin et al., 2009) is then performed to obtain apparent phase velocities using phase travel times over 0.05°-by-0.05° grid points for each 182 virtual source and period. The quality of the derived phase velocity is improved by stacking

183 phase velocities obtained from all available virtual sources. The resulting isotropic phase

velocities increase with depth and are generally higher in mountainous regions, such as the

185 Peninsular Ranges, and lower in basins, including the LA Basin, Salton Trough and Central

Valley (Fig. 6a,c). The results are generally consistent with corresponding phase velocity

measurements in the area based on beamforming analysis (Roux & Ben-Zion, 2017). Phase
 velocity uncertainty is determined based on the variation of measurements using different virtual

- 189 sources (Fig. 6b,d).
- 190 191

2.5 Monte Carlo Joint Inversion

To jointly invert phase velocities and H/V ratios for shear wave velocity, we use a nonlinear Bayesian Markov Chain Monte Carlo (MCMC) method. This method has several advantages: it fully explores the available parameter space, it is unlikely to be trapped in a local minimum, and can fully quantify model uncertainty (Shen et al., 2012; Roy and Romanowicz, 2017).

Following Shen et al. (2012), we assemble a starting model across Southern California 197 198 with uniform horizontal 0.05°-by-0.05° spaced grid points based on the SCEC Community Velocity Model CVMS4.26 (Lee et al., 2014). At each grid point, we extract an independent 1D 199 Vs model which we parameterize with three layers: a linear sedimentary layer near the surface, a 200 201 crustal layer described by 10 cubic B-splines and the upper mantle defined by 5 cubic B-splines to a total depth of 50km. Our solution is inherently regularized due to the spatial discretization of 202 the B-splines employed, which results in uncertainty of our final model to be underestimated 203 (Dettmer et al., 2016). The thickness of the top linear layer is defined by the depth which the 204 CVMS4.26 model reaches 2.3 km/s, roughly representing a sedimentary layer with linearly 205 increasing Vs. The depth to the Moho is explicitly defined in the CVMS4.26. We choose to use 206 10 cubic B-splines in the crust to honor the often-presented complexity in the starting reference 207 model. To decrease the number of parameters in our MCMC inversion, we only perturb the even 208 number of the crustal spline value and using the mean of the neighboring spline perturbation to 209 determine the odd number of the spline perturbation. This also allows our model to honor the 210 basic structure resolved in the CVMS4.26 model (Lee et al., 2014). We hold the Moho depth and 211 mantle splines completely fixed. Holding the mantle parameters fixed is a reasonable approach 212 because the H/V and phase data at the periods we employ have very weak sensitivity below the 213 214 crust (Fig. 3). In this study, we use the empirical relationships described by Brocher (2005) to determine Vp and density from the Vs models. 215

For each 1D model, we next create *a priori* distributions of the 8 free parameters: the Vs velocity values of the top and bottom of the sedimentary layer, sedimentary thickness, and 5 spline values in the crust. The prior distributions are centered around the starting model values and are obtained by the Gaussian probability distribution described in Table 1, where the Gaussian width of each parameter is empirically chosen to provide full sampling of the model space (Shen et al., 2012). These *a priori* distributions control the parameter space explored by the inversion.

We follow the Markov Chain Monte Carlo inversion described by Shen et al. (2012) to determine the posterior distribution from the prior distributions. Models are randomly selected from the prior distributions by simultaneously perturbing all eight parameters. The model misfit is then characterized as the χ^2 difference between the observed and the forward-calculated H/V and phase velocity of each model using the method of Herrmann et al. (2004). The χ^2 model

misfit treats phase velocity and H/V measurements equally, where phase velocity and H/V228 229 uncertainty corresponds to 150% of the standard deviation of the mean for the associated measurement at each period. We use 150% of the standard deviation to account for potential 230 systematic bias that is not encompassed by the measurement variation (Lin et al., 2012; Shen et 231 al., 2013). Parameter space is then explored following the Metropolis algorithm (e.g., Mosegaard 232 & Tarantola, 1995; Beichl et al., 2000). We obtain the likelihood functional of the model from 233 forward computation using the Thomson-Haskell method, computed via Herrmann et al. (2004), 234 with an earth-flattening transformation (Shen et al., 2012). If the probability of acceptance, 235 related to the misfit of the model to the data through the likelihood function (Shen et al., 2012), 236 is higher than the previous model we define a new perturbation from this model. If the model has 237 poorer misfit we either instead define a new perturbation from the previous model or accept this 238 realization. This decision is guided by a probability defined by the likelihood function and 239 Metropolis law, as discussed in Shen et al. (2012), and prevents the inversion process from 240 becoming trapped in a local minimum (Mosegaard & Tarantola, 1995). After 3000 iterations of 241 random perturbations, we begin a new random set of model iterations from the original starting 242 model. We perform 10 of these jumps with 3000 iterations per jump. To form the posterior 243 244 distribution, we select all models with misfit less than or equal to 1.5 times the misfit of the absolute lowest misfit model. This estimation of the posterior distribution is computationally 245 efficient and effectively removes models that may have been accepted during the inversion prior 246 247 to the equilibrium state, which describes the posterior model distribution. However, by doing so we also effectively trim our posterior distribution and underestimate true posterior distribution 248 width. The model created by the mean of each parameter of models within the posterior 249 distribution is our final model. To obtain a full 3D model across Southern California, we perform 250 this inversion independently for each 0.05°x0.05° grid point. 251

To increase inversion stability, we add a few reasonable constraints: the maximum Vs in 252 the crust is 4.9 km/s, the sedimentary layer must have increasing Vs with depth, the first two 253 splines in the crust must be increasing, and a positive change in velocity must exist across the 254 (top sedimentary) linear layer to the (middle) crustal layer. The mean number of accepted 255 posterior models for each 1-dimensional inversion is 136 models. Example MCMC inversions 256 for two grid points are shown in Figures 7 and 8 (locations are shown with a green and blue star 257 each in Figure 1). This includes the starting model formed from the CVMS, the entire model 258 space searched, posterior models, the final mean result and associated standard deviation 259 (Figures 7a & 8a). Additionally, the forward model results of H/V and phase velocity are shown 260 for the data, the starting model, all posterior models and the final mean model (Figures 7b,c & 261 8b,c). Finally, the distributions of posterior compared to prior parameters are shown for several 262 different parameters (Figures 7d,e,f & 8d,e,f). From these distributions it is evident that our final 263 results show a narrow Gaussian distribution and are sensitive to the shallow and mid-crust shear 264 wave velocity. 265

266 **3 Results**

267

3.1 H/V and Phase Velocity Results

The map-view images of H/V and phase velocity (Figures 5a,c and 6a,c) show consistent patterns related to geologic structure. Regions of high H/V and low phase velocity correspond to sedimentary basins including the LA Basin, Central Valley and Salton Trough, and the size of these features decreases with increasing period (corresponding to depth). Regions with low H/V and high phase velocity correspond to mountainous regions, including the Sierra Nevada and

Peninsular Ranges. Major faults including the San Andreas, San Jacinto, and Garlock faults, 273 274 appear as sharp boundaries separating regions of different velocity and H/V.

The corresponding uncertainty maps (Figures. 5b,d and 6b,d) provide confidence in the 275 derived phase velocity and H/V data. We see a decrease of uncertainty with increase in period 276 for H/V and phase velocity, with uncertainties less than 5% at all periods. Specifically, we 277 observe variations in uncertainty up to 5% with most uncertainties less than 2% for phase 278 velocity, but slightly higher variations in H/V uncertainties, up to 7% at shorter periods, with 279 most uncertainties less than 5%. Higher uncertainties over lower periods is likely due to stronger 280 heterogeneity and wavefield complexity, such as multipathing and off-great-circle propagation, 281 in the shallow structure. Both datasets have relatively higher uncertainty near the edges of the 282 region, which may be attributed to poorer station distribution and azimuthal coverage. 283 Observations of high uncertainties at low periods for both datasets in basin areas (e.g., Salton 284 Trough, LA Basin, Central Valley) may be due to strong 3D heterogeneity effects in the shallow 285 structure. In order to account for underestimation of uncertainty in our measurements, we use 286 150% of the measured uncertainty for both phase velocity and H/V in the joint inversion. 287

288

3.2 Monte Carlo Inversion 1-D Results

289 Figures 7 and 8 summarize the 1D inversion results for two different locations: the San 290 Gorgonio Pass and the Central Valley respectively. By using a starting model that predicts the 291 292 phase velocity dispersion fairly accurately, we are able to search the full model space and find a suite of best-fitting models to form the posterior. In both example cases, the starting model 293 performs poorly at predicting H/V ratio. This is because the CVMS4.26 model (Lee et al., 2014) 294 was developed using data with limited sensitivity to the shallow structure (e.g. the top 3 km). 295 The inclusion of H/V data in the present work leads to strong changes in the shallow structure 296 (<10km depth), which dramatically improves H/V fit in addition to slightly improving phase 297 velocity fit. In nearly all cases, changes related to higher H/V in the data than the starting model 298 correspond to low velocity zones in the upper few km that are completely absent in the starting 299 model, especially near the Salton Trough, Coast Ranges and Indian Wells Valley, as discussed in 300 Section 4. Because of the shallow sensitivity of the datasets, we do not constrain structure below 301 \sim 25km depth; the posterior distributions are quite broad and simply average back to the starting 302 model below this depth. Comparing prior and posterior model distributions (Figures 7d,e,f and 303 8d,e,f) indicates that the posterior models lie completely within the examined parameter space, 304 are much more tightly constrained by the data than the prior distributions, and generally follow 305 Gaussian distributions. In addition, the posterior distributions become wider with depth, 306 indicating a relative decrease in model certainty as expected. 307

308 309

3.3 Shear Velocity in 3D

310 The joint inversion results in terms of both absolute velocity and relative change to the starting model CVMS4.26 are presented in Figures 9 and 11, and the distribution of misfit values 311 312 is displayed in Figure 10. We also include corresponding plots in the Supporting Information Section (Figures S1-S3) that compare our joint inversion results to the SCEC Community 313 Velocity Model – Harvard. Specifically, map-view slices at several depths are shown in Figure 9, 314 and seven cross-sections are shown in Figure 11. The map-view images have no smoothing 315 applied after the individual 1D inversions, while the cross-sections were created using narrow-316 width cubic interpolation to sample along the arbitrarily-oriented profiles (Figure 9e). The final 317 318 model shows the strongest changes from the starting model at shallow depths due to the addition of H/V ratios in the inversion. These H/V data in general require lower velocities near the 319

- surface (<2km depth) and higher velocities in the upper crust compared to the starting model.
- 321 The correspondence of the absolute velocity structure to various geological provinces is
- discussed in detail in Section 4.
- The total χ^2 misfit at each grid point for both the starting and final models is shown in
- map view for each dataset (Figure 10). The final model improves fit of both phase velocity and
- H/V compared to the starting model. In general, misfit of the final model is low (< 1.5), except for the Los Angeles basin and in a few localized areas at the edges of the imaged region. This
- relatively high misfit in the LA basin is potentially due to the Moho being fixed at an incorrect
- depth in the present work; results from a recent dense seismic array indicate that the Moho
- beneath the LA basin is much shallower than previously thought (Ma & Clayton, 2016).
- Nevertheless, the new joint inversion model features significantly better fits to the data than the
- 331 CVM-S starting model, including in the LA basin.

332 4 Discussion

We provide interpretations of the 3D joint inversion model in the context of the various geological provinces of Southern California. Though a comprehensive interpretation of a seismic velocity model in terms of geology requires consideration of a variety of additional parameters such as temperature, fluid content, fracture density, and Poisson's ratio (e.g., Christensen, 1996; Karato & Jung, 1998), we note the strong correspondence of the present model to previous results and to expectations based on geological inferences.

339 340

4.1 Southern Central Valley and Sierra Nevada

At 0.5 and 2 km depth (Figure 9a,b), a prominent feature in our result is the transition 341 from the slower sediments in the Central Valley to the faster Sierra Nevada foothills. This 342 transition is consistent with previous tomographic imaging results (Tape et al., 2010; Lee et al., 343 2014). However, the improved shallow sensitivity provided by the H/V data leads to a reduced 344 sediment thickness (3-4km deep) in the southern tip of the Central Valley (F-F') which more 345 closely matches the active-source studies (Fliedner et al., 2000) and focused ambient noise 346 imaging studies (Fletcher & Erdem, 2017). As evident in Figure 9(a), to the east of the Sierra 347 Nevada range and north of the Garlock fault is a low-velocity zone in the Indian Wells Valley, 348 which has previously been imaged as a shallow (<3km) feature (Tape et al., 2010) that may host 349 350 enhanced hydrothermal activity (Ho-Liu et al., 1988).

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- 352 353

4.2 Coast Ranges and Transverse Ranges

The Coast Ranges, west of the Central Valley, are very slow with similar velocities to the 354 Central Valley (Figure 9). This has been observed in previous studies (Tape et al., 2010) and 355 attributed to stacking of Miocene and younger sediments caused by East-West shortening along 356 various active structures (Namson & Davis, 1988). At 0.5 km depth, the Transverse Ranges and 357 358 surrounding areas have blocks of varying slow velocities corresponding to major faults including the San Gabriel fault, San Cayetano fault and San Andreas fault (A-A'). The southern San 359 Andreas fault has clear velocity contrasts with opposite polarities to the NW and SE of San 360 Gorgonio. These features were observed with fault zone head waves and previous tomographic 361 results and can have important implications for the size and directivity of large earthquakes on 362 the southern San Andreas fault (Share & Ben-Zion, 2016). Evidence of fold and thrust belts 363 corresponding to the Ventura anticline and San Cayetano fault can be seen to the southwest of 364

the San Cayetano fault surface trace, and supports the model proposed by Hubbard et al. (2014).
The San Cayetano fault also bounds a fast 5km feature (F-F') and marks the eastern edge of
slow, shallow sediments which transition to much faster crystalline rocks (Powell, 1981)
bounded to the cast by the San Andreas fault

bounded to the east by the San Andreas fault.

4.3 Mojave Desert

The Mojave Desert is bounded by the Garlock fault to the north and the San Andreas fault to the southwest (Fig. 1). South of the Sierra Nevada range, the Garlock fault separates the faster mountains from slower material in Antelope Valley in the westernmost edge of the Mojave, as seen at very shallow depths (Figure 9a). Contrastingly, at 2km depth, Antelope Valley is faster than the Sierra Nevada range, north of the Garlock fault, or the region south of the San Andreas fault. Previous active-source studies (Lutter et al., 2004) have also shown that this shallow valley (G-G') contains slow material overlying fast material.

This trend is also seen in the CVMS, but is deeper than the new model (Figure 9 d.e). 378 Mid-crustal structure beneath the Mojave Desert is known to be strongly anisotropic (e.g., Louie 379 & Clayton, 1987) due to complicated Miocene metamorphic processes (Fletcher et al., 1995). 380 Additionally, velocity contrasts are seen across the Eastern California Shear Zone similar to 381 previous results (Tape et al., 2010; Lee et al., 2014). The northeastern part of the Mojave (A-A') 382 has slower material near the surface overlying the shallow fast material beginning at the 383 Lockhart fault. At 9km depth, the region south of the Garlock fault appears slower than the 384 385 material in the Sierra Nevada mountains north of the Garlock fault, similar to previous results (e.g., Tape et al., 2010). 386 387

4.4 Los Angeles Basin

Los Angeles is underlain by a well-studied basin with a mean depth of ~5km and 389 maximum depth of 10km (Magistrale et al., 1996) that has been developing since at least the 390 Middle Miocene (Ingersoll & Rumelhart, 1999). The internal structure of the basin is 391 complicated by multiple east-west trending active blind thrust faults (Shaw et al., 2015) with 392 393 complex patterns of interaction (Rollins, 2018). The LA Basin in the new model has the lowest velocities in the entire region, though it is also the region with the highest residual misfit (Fig. 394 10) likely due to the incorrectly constrained Moho (Ma & Clayton, 2016) as described in Section 395 3.3. The overall basin structure (B-B' and G-G') is generally consistent with CVMS-4.26, 396 though with much lower velocities at depths less than 2km as required by the H/V data. These 397 lower velocities are crucial to correctly quantify seismic hazard (e.g., Olsen, 2000) in the LA 398 399 Basin, the most populous region in Southern California with some of the highest seismic risk in the United States (e.g., Petersen et al., 2015). 400

There are clear velocity changes across the Newport-Inglewood fault (B-B', 1-3km 401 depth), which have previously been observed in more detailed studies (Lin et al., 2013). The 402 basin deepens between the Newport-Inglewood and Whitter faults, in line with previous 403 geological (Shaw & Suppe, 1996) and geochemical (Boles et al., 2015) studies. Additionally, 404 405 there is a strong contrast across the Whitter fault leading to a shallower basin NE of the fault as expected from the fault throw (Davis et al., 1989; Shaw et al., 2015). The basin is bounded 406 sharply to the North by the Sierra Madre fault (B-B', G-G'), an active reverse fault largely 407 responsible for the uplift of the San Gabriel Mountains (Shen et al., 2011). 408

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4.5 San Bernardino Basin and Major Faults

The San Bernardino basin is a shallow (<2km; Anderson et al., 2004) feature bounded to 411 the east by the San Andreas Fault and to the west by the San Jacinto Fault zone (Fig 9a). To the 412 southeast of the basin is a region of small-scale faults likely responsible for transfer of stress 413 from the southern San Andreas fault to the more favorably oriented San Jacinto fault 414 (Langenheim et al., 2004; Fialko, 2006). The San Jacinto bounds some of the strongest across-415 fault contrasts in the entire model, as seen in previous tomographic studies (Tape et al., 2010; 416 Allam & Ben-Zion, 2012; Allam et al., 2014). This relatively fast region is composed of 417 Cretaceous plutons (Morton & Kennedy, 2005) and is also sharply bounded to the SW by the 418 Elsinore fault (C-C', D-D'). The sense of the velocity contrast changes from SW-fast to NE-fast 419 southward along the fault due to the presence of the large San Jacinto plutons (Hill, 1988). The 420 421 sharp across-fault velocity contrasts can lead to a preference for NW-propagating earthquake ruptures (Shi & Ben-Zion, 2006; Allam et al., 2014), which can be up to M7.5 based on 422 paleoseismic data (Rockwell et al., 2015). 423

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4.6 Salton Trough and Peninsular Ranges

The Peninsular Ranges in far southwestern California are composed of a series of large-426 scale Mesozoic plutonic rocks (Gastil, 1975; DePaolo, 1981) that are almost completely 427 unfaulted (Plesch et al., 2007). Bounded to the east by the Elsinore fault, the Peninsular Ranges 428 have the fastest shallow velocities in the present model (Figure 9a), in agreement with previous 429 regional tomographic studies (Tape et al., 2010; Lee et al., 2014). To the east is the Salton 430 Trough, a region of crustal extension (Sylvester & Smith, 1976) with an extremely shallow 431 Moho (Ozakin & Ben-Zion, 2015) possibly indicating incipient mid-ocean spreading centers 432 (Robinson et al., 1972; Han et al., 2016). The Salton Trough in the present model is much wider 433 and slower in the upper few km (Figure 9a; Figure 11 E-E') than the CVMS4.26, in agreement 434 with previous active-source studies (Livers et al., 2012; Fuis et al., 2017; Han et al., 2016). The 435 Superstition Hills fault and Brawley Seismic zone have no obvious signal in the seismic velocity 436 model, supporting the idea that these are regions of distributed deformation due to multiple fault 437 438 strands (Hudnut et al., 1989) and diffuse seismicity (Geng et al., 2013).

439 **5** Conclusions

We combine Rayleigh-wave H/V ratios and phase velocity measurements in a joint 440 Bayesian inversion to determine a regional shear velocity model for Southern California with 441 improved resolution in the surface, shallow and upper crustal structure. Previous models such as 442 the CVMS4.26 (Lee et al., 2014) have incorporated information from ambient noise and full 443 waveforms but did not incorporate amplitude information and therefore have a relatively weak 444 constraint on structure above 3km depth. By combining H/V ratios and phase velocity 445 measurements, we gain sensitivity to shallow and mid-crustal shear velocity structure. The 446 obtained large-scale mid-crustal features are similar to previous high-resolution models (e.g., 447 Tape et al., 2010; Lee et al., 2014; Barak et al., 2015; Fang et al., 2016), lending confidence in 448 the new model overall. The main improvement is the addition of new shallow features in the 449 updated model, including more accurate basin depths and other near-surface low-velocity zones 450 that have strong implications for studies of seismic hazard. The final model is a self-consistent 451 regional-scale seismic velocity model with resolution from seismogenic depth to the surface. 452 In addition to resolving large-scale features of the crust, our shear velocity model 453 includes small-scale shallow structure previously only seen by local imaging studies (Allam et 454 al., 2014; Fliedner et al., 2000; Lin et al., 2013; Fuis et al., 2017). In the north this includes the 455

- 456 shallower-sediments in the southern tip of the Central Valley (Fliedner et al., 2000; Fletcher &
- 457 Erdem, 2017), high velocity of the Sierra Nevada Range (Tape et al., 2010), shallow slow
- velocities in the Coast and Transverse Ranges (Tape et al., 2010) and evidence of fold and thrust
- 459 faults (Hubbard et al., 2014). We resolve similar shallow structure in the LA basin to the
- 460 CVMS4.26 (Lee et al., 2014) while also imaging the Newport-Inglewood fault (Lin et al., 2013)
- and Whittier faults (Shaw and Suppe, 1996). We also are able to see the San Bernardino basin
 and differing velocity structure across the Elsinore, San Jacinto and San Andreas faults (Allam et
- and differing velocity structure across the Elsinore, San Jacinto and San Andreas faults (Allam et al., 2014; Allam and Ben-Zion, 2012; Zigone et al., 2015). In the southern end of the region, we
- al., 2014; Allam and Ben-Zion, 2012; Zigone et al., 2015). In the southern end of the region, we
 recover the Salton Trough and Peninsular Range with similar structure to active source studies
- 464 (Livers et al., 2012; Fuis et al., 2017; Han et al., 2016). Our results demonstrate the considerable
- improvement to ambient noise imaging that can be gained from the incorporation of spatially
- 467 dense Rayleigh wave H/V measurements to constrain shallow structure.

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- 484 485

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817 **Table 1**

818 Prior Distributions in Joint Inversion

Parameters	Range	Gaussian width
Sedimentary thickness	$10 \pm m_0 (km)$	0.2 km
Vsv, top of sedimentary layer	$m_0 \pm 0.5 m_0 (\text{km s}^{-1})$	0.1 km s ⁻¹
Vsv, bottom of sedimentary	$m_0 \pm 0.5 m_0 (\text{km s}^{-1})$	0.1 km s ⁻¹
layer		
Crust 0 th B-spline	$m_0 \pm 0.5 m_0 (\text{km s}^{-1})$	0.2 km s^{-1}
Crust 2 nd B-spline	$m_0 \pm 0.4 m_0 (\text{km s}^{-1})$	0.2 km s^{-1}
Crust 4 th B-spline	$m_0 \pm 0.4 m_0 (\text{km s}^{-1})$	0.2 km s^{-1}
Crust 6 th B-spline	$m_0 \pm 0.3 m_0 (\text{km s}^{-1})$	0.2 km s^{-1}
Crust 8 th B-spline	$m_0 \pm 0.2 m_0 (\text{km s}^{-1})$	0.2 km s^{-1}

819 (left) A full list of the inversion parameters, (middle) the ranges explored, and (right) the

gaussian half-width used to define the *a priori* distributions

Figure 1. Location map of the imaged region. Stations (blue triangles), faults (black lines), topography (grayscale), and various sub-regions are shown. Example stations FMP (yellow

- triangle) and WBS (red triangle) are marked, with corresponding cross-correlation H/V
- measurement distributions shown in Figures 2 & 4 respectively. Green and blue stars mark
- locations of example joint inversion results in Figures 7 and 8. Major geological features
- mentioned in the text are labeled in full, with the following abbreviations for major faults: San
- Andreas (SAF), Garlock (GF), Elsinore (EF) and San Jacinto (SJF).
- 829

Figure 2. (a) Four-component (ZR, ZZ, RR, RZ) ambient noise cross-correlations between
stations FMP and WBS bandpassed around 8 s period. Clear Rayleigh waves are visible on all
causal components. (b) The 8-s Rayleigh-wave particle motion in radial and vertical directions
observed at receiver station WBS (left) excited by a vertical force (top) or horizontal force
(bottom) at station FMP. (Right) same as (left), but with FMP being the receiver and WBS the
virtual source. Stations locations are shown in Figure 1.

- 836
- Figure 3. (a) H/V and (b) phase velocity sensitivity kernels for a location near San Jacinto fault
 (Figure 1 green star) at three different periods based on the CVMS4.26 shear-wave velocity
 model.
- 840

Figure 4. Distributions of H/V measurements from individual cross-correlations found from the station acting as a source (blue) or receiver (magenta) with stabilization range marked with gray dashed lines. Measurement distributions for station FMP at (a) 7 s period and (c) 15 s period. Measurement distributions for station WBS at (b) 7 s period and (d) 15 s period.

- 845
- **Figure 5.** H/V and uncertainty maps showing individual station results (circles) and
- corresponding interpolation over the entire region. (a) H/V measurements at 7 seconds period.
- (b) H/V uncertainty, shown as the ratio of standard deviation of the mean to H/V value, for 7 s
- period. (c) Similar to (a) but for 15 s period. (d) Similar to (b) but for 15 s period.
- 850

Figure 6. Phase velocity and uncertainty maps. (a) Phase velocities (Vph) at 7 seconds period.
(b) Vph uncertainty, shown as the ratio of standard deviation of the mean to Vph value, for 7 s
period. (c) Similar to (a) but for 15 s period. (d) Similar to (b) but for 15 s period.

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Figure 7. Example 1D joint inversion result at a location near the San Andreas fault (Figure 1 855 green star). (a) Shear wave velocity versus depth showing initial model as red triangles (CVMS), 856 full range of model space searched (green dashes), posterior models (cyan lines), the final model 857 (white dots), and final model standard deviation (black lines). The fixed Moho depth is also 858 marked in gray. (b) H/V dispersion curves including H/V data (black dots), 150% of the 859 860 uncertainty (error bars), predicted H/V ratios obtained from the starting model (red triangles), posterior models (cyan lines), and the final average model (white dots). (c) Similar to (b), but for 861 phase velocities. (d) The *a priori* distribution of Vs at 0.5km depth shown in transparent 862 histogram with thick black lines and posterior distribution in blue. The mean posterior parameter 863 and standard deviation are shown. (e) Similar to (d) but for 2 km depth (f) Similar to (d) but for 9 864 km depth. 865

Figure 8. Same as Figure 7 but for a station in the Central Valley (Figure 1 blue star). Note the stronger gradient and shallower sediment depth in the final model (white dots) compared to the starting model (red triangles) in (a). The fit to the phase data is slightly improved (c), but the H/V fits are dramatically improved over the starting model (b).

- 871
- Figure 9. Joint inversion shear velocity (Vsv, km s⁻¹) results for depth of (a) 0.5 km, (b) 2 km, and (c) 9 km. Also shown are the differences between final and starting models for depths of (d)
- 0.5 km (e) 2 km (f) 9 km. Cross-sections denoted in (e) are shown in Figure 11.
- 875

Figure 10. χ^2 misfit over all periods to H/V and phase velocity from the final model for (a) H/V, (b) phase velocity, and (c) joint H/V and phase velocity. Shown in (d)-(f) are results similar to (a)-(c), but from the starting model (CVMS). Note the wider scale bar for the starting compared to the final model; misfit is lower for all datasets in the final model.

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Figure 11. Cross-Sections (Figure 9e) of (left) final inversion Vsv results and (right) difference

between final and initial (CVMS) Vsv. (a) A-A' cross section with the San Cayetano, San
Gabriel, Clearwater, San Andreas, Lockhart and Garlock Fault surface traces marked. (b) B-B'

Gabriel, Clearwater, San Andreas, Lockhart and Garlock Fault surface traces marked. (b) B-B'
 cross-sections with the Newport-Inglewood, Whittier, Sierra Madre and San Andreas fault

surface traces marked. (c) C-C' cross-sections with Elsinore, San Jacinto, Banning and Mill

Creek fault surface traces marked. (d) D-D' cross-section with Elsinore, San Jacinto, Baining and With Creek fault surface traces marked. (d) D-D' cross-section with Elsinore, San Jacinto and San

Andreas fault surface traces marked. (e) E-E' cross-section with Elsinore, San Jacinto and San 887 Andreas fault surface traces marked. (e) E-E' cross-section with Elsinore and Superstition Hills

faults and Brawley seismic zone labeled. (f) F-F' cross-section with San Cayetano and San

Andreas faults marked. (g) G-G' cross section with Newport-Inglewood, Sierra Madre, San

890 Gabriel, San Andreas, and Garlock faults marked.

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Figure 1.



Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.



Figure 7.



Figure 8.



Figure 9.



Figure 10.



Figure 11.

